

Dynamic Bifurcation from Infinity of Nonlinear Evolution Equations*

Chunqiu Li[†], Desheng Li[‡], and Zhijun Zhang[§]

Abstract. This paper is concerned with dynamic bifurcation from infinity and multiplicity of stationary solutions for nonlinear evolution equations near resonance. First, we prove some new global continuation results and establish a general theorem on dynamic bifurcation from infinity in the framework of local semiflows on metric spaces. Then, by applying these abstract results, we derive more precise descriptions on the dynamic bifurcation from infinity of evolution equations in Banach spaces. Finally, we focus our attention on the parabolic equation $u_t - \Delta u = \lambda u + f(x, u)$ associated with the Dirichlet boundary condition, where f satisfies appropriate Landesman-Laser type condition. A detailed discussion on the dynamical behavior and the multiplicity of stationary solutions of the equation near resonance will be presented.

Key words. Conley index, nonlinear evolution equation, bifurcation from infinity, parabolic equation, resonance

AMS subject classifications. 37B30, 35B32, 35K55, 34C23

1. Introduction. This paper is concerned with the nonlinear evolution equation

$$(1) \quad \frac{du}{dt} + Au = \lambda u + f(u, \lambda)$$

on a Banach space X , where A is a sectorial operator on X with compact resolvent, $\lambda \in \mathbb{R}$ is the bifurcation parameter, and $f(u, \lambda)$ is a locally Lipschitz continuous mapping from $X^\alpha \times \mathbb{R}$ ($0 \leq \alpha < 1$) to X which is sublinear as $\|u\|_\alpha \rightarrow \infty$ uniformly on bounded λ -intervals. We are basically interested in the dynamic bifurcation from infinity of the equation and its applications.

This topic can be traced back to the work of Rabinowitz [30], in which the author studied the bifurcation from infinity of stationary solutions of the equation in a general setting of operator equations of the following form:

$$(2) \quad u = \lambda Lu + K(u, \lambda),$$

where L is a compact operator, $\lambda \in \mathbb{R}$, and $K(u, \lambda) = o(\|u\|)$ as $\|u\| \rightarrow \infty$ uniformly on bounded λ -intervals. It was shown that if μ^{-1} is a real eigenvalue of L of odd multiplicity, then (∞, μ) is a bifurcation point. Furthermore, there is a continuum of solutions of (2) which goes to infinity as $\lambda \rightarrow \mu$. This result was partially extended by Toland [38], Dias and Hernandez [10] and Schmitt and Wang [36] to potential operator equations to cover the case of even multiplicity. The interested reader is referred to [1, 2, 5, 6, 13, 16, 22, 23, 27, 28, 31, 33],

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[†]School of Mathematics, Tianjin University, Tianjin 300072, China (*Email:* licqmath@tju.edu.cn).

[‡]School of Mathematics, Tianjin University, Tianjin 300072, China; Center for Applied Mathematics, Tianjin University, Tianjin 300072, China (*Corresponding author. *Email:* lidsmath@tju.edu.cn, lidsmath@163.com).

[§]School of Mathematics, Yantai University, Yantai 264005, China (*Email:* zjzhang@ytu.edu.cn).

32 etc. for concrete examples on the bifurcation from infinity and multiplicity near resonance for
 33 differential equations under various boundary conditions.

34 For a nonlinear system as in (1), stationary solutions may be far from being adequate
 35 for understanding its dynamics. This is because that the dynamics of the system is usually
 36 determined not only by its stationary solutions, but also by all other bounded full ones. In
 37 fact, it is often the case that a system may have no stationary solutions. It is therefore of
 38 great importance to develop appropriate theories to analyze the bifurcation of bounded full
 39 solutions. A fundamental one in this line is the well-known Hopf's bifurcation theory, which
 40 was first developed in the very early work of Poincaré [29] around 1892. Actually it forms the
 41 central part of the classical dynamic bifurcation theory. The Hopf's bifurcation theory focuses
 42 on the case when there are exactly a pair of conjugate eigenvalues of the linearized equation
 43 crossing the imaginary axis, and was fully developed in the 20-th century. One can find a vast
 44 body of literature on how to determine Hopf bifurcations for nonlinear systems arising from
 45 applications. To deal with the general case, some other dynamic bifurcation theories need to
 46 be developed, and the Conley index theory, attractor theory and so on allow us to take a step;
 47 see e.g. [12, 18, 21, 34, 39, 40], etc.

48 Very recently, Li and Wang [18] established some new local and global bifurcation results
 49 in terms of invariant sets via the Conley index theory, completely extending the well-known
 50 Rabinowitz's global bifurcation theorem to the dynamic bifurcation of nonlinear evolution
 51 equations without requiring the "crossing odd-multiplicity" condition. Inspired by this work
 52 and some other ones mentioned above, in this paper we consider the dynamic bifurcation
 53 from infinity of (1). This problem was actually addressed in Ward [40]. The author first
 54 established a global continuation theorem (see Remark 3.2). Then he proved the following
 55 interesting result: For any real numbers $c < d$ such that the interval $[c, d]$ contains exactly
 56 one number $\mu \in \operatorname{Re} \sigma(A)$ with $c < \mu < d$, there exists a continuum $\mathcal{C} \subset X^\alpha \times [c, d]$ meeting
 57 $X^\alpha \times \{c, d\}$ such that

(1) for $\lambda \neq c, d$, $\mathcal{C}[\lambda]$ consists of bounded full solutions, where

$$\mathcal{C}[\lambda] = \{x : (x, \lambda) \in \mathcal{C}\};$$

58 (2) there is a sequence $\lambda_n \rightarrow \mu$ such that $\mathcal{C}[\lambda_n]$ is unbounded as $\lambda_n \rightarrow \mu$.

59 (See [40, Theorem 3.2].) Note that the continuum \mathcal{C} in the above result may contain either
 60 all the connected branches of bounded full solutions of the equation meeting $X^\alpha \times \{c\}$, or all
 61 the connected branches of bounded full solutions meeting $X^\alpha \times \{d\}$, according to which side
 62 \mathcal{C} will meet. Here, by using the techniques in [18] we will prove some new continuation results
 63 and establish an abstract theorem on bifurcation from infinity in terms of local semiflows on
 64 metric spaces. Then based on these theoretical results, we give some more precise descriptions
 65 on the dynamic bifurcation from infinity for (1).

66 As an example, we consider the parabolic equation

67 (3)
$$u_t - \Delta u = \lambda u + f(x, u), \quad x \in \Omega$$

68 associated with the homogeneous Dirichlet boundary condition, where Ω is a bounded domain
 69 in \mathbb{R}^n , and f is a bounded function satisfying the following Landesman-Laser type condition:

70

$$71 \quad (4) \quad \liminf_{s \rightarrow +\infty} f(x, s) \geq \bar{f} > 0, \quad \limsup_{s \rightarrow -\infty} f(x, s) \leq -\underline{f} < 0$$

uniformly for $x \in \bar{\Omega}$ (where \bar{f} and \underline{f} are independent of x). First, we give a detailed discussion on the dynamic bifurcation from infinity of the equation near any eigenvalue μ_k of the operator $-\Delta$ (in $H_0^1(\Omega)$). Specifically, we prove that there exists $\delta > 0$ such that for each $\lambda \in \Lambda_- = [\mu_k - \delta, \mu_k)$, the maximal compact invariant set S_λ of the equation has a Morse decomposition $\mathcal{M} = \{M_\lambda^\infty, M_\lambda^1\}$ with M_λ^1 being uniformly bounded on Λ_- while

$$\lim_{\lambda \rightarrow \mu_k^-} \min_{v \in M_\lambda^\infty} \|v\| = \infty.$$

Besides, there is at least one connecting trajectory γ between M_λ^∞ and M_λ^1 . More interestingly, it will be shown that each of the following two sets

$$\mathcal{K}^1 = \overline{\bigcup_{\lambda \in \Lambda_-} (M_\lambda^1 \times \{\lambda\})}, \quad \mathcal{K}^\infty = \overline{\bigcup_{\lambda \in \Lambda_-} (M_\lambda^\infty \times \{\lambda\})}$$

contains a connected component Γ with

$$\Gamma[\lambda] := \{u : (u, \lambda) \in \Gamma\} \neq \emptyset, \quad \forall \lambda \in \Lambda_-.$$

72 The bifurcation and multiplicity of elliptic equations near resonance is always an inter-
 73 esting topic and has attracted much attention in the past decades. As a byproduct of our
 74 dynamical argument, we can naturally derive some bifurcation and multiplicity results on the
 75 stationary problem:

$$76 \quad (5) \quad \begin{cases} -\Delta u = \lambda u + f(x, u), & x \in \Omega, \\ u(x) = 0, & x \in \partial\Omega. \end{cases}$$

77 This problem was first studied by Mawhin and Schmitt [22], where the authors considered the
 78 case when λ crosses an eigenvalue of odd multiplicity. Later Schmitt and Wang [36] developed
 79 a theory on bifurcation from infinity for potential operators, through which they extended
 80 the results in [22] to the case when λ crosses any eigenvalue μ_k . More specifically, under
 81 an abstract Landesman-Lazer type condition on the Nemitski operator $\tilde{f} : H_0^1(\Omega) \rightarrow L^2(\Omega)$
 82 corresponding to the function $f(x, s)$, the authors proved the following result and its “dual”
 83 version: there exists $\delta > 0$ such that for each $\lambda \in (\mu_k, \mu_k + \delta]$ the equation (5) has at least
 84 two solutions with one of which approaching ∞ as $\lambda \rightarrow \mu_k$, and for each $\lambda \in [\mu_k - \delta, \mu_k]$ it has
 85 at least one. (Some further development and extensions can be found in [4, 6, 9, 11, 27, 37],
 86 etc.) As an application of our dynamical bifurcation results, we show that the “dual” version
 87 of the above result holds true under the hypothesis (4). Moreover, there exists an open dense
 88 subset \mathcal{D} of \mathbb{R} such that for $\lambda \in \Lambda_- \cap \mathcal{D}$, where $\Lambda_- = [\mu_k - \delta, \mu_k)$, the problem has at least
 89 three distinct solutions.

90 Special attention will also be paid to the case where $f(x, s) = o(|s|)$ as $|s| \rightarrow 0$ uniformly
 91 for $x \in \bar{\Omega}$, in which we can say a little more on the multiplicity of nontrivial solutions of (5).
 92 Such a case was studied in Chiappinelli, Mawhin and Nugari [6], where the authors considered

93 the multiplicity of solutions of the problem near the first eigenvalue μ_1 . Under appropriate
 94 Landesman-Laser type conditions, it was proved, among other things, that the problem has
 95 at least two distinct nontrivial solutions as $\lambda \rightarrow \mu_1^+$. (We mention that the nonlinearity
 96 in [6] was allowed to be unbounded.) Here we present some more precise information on
 97 the multiplicity of solutions for the problem near any eigenvalue μ_k under the condition (4).
 98 Roughly speaking, we show that (5) has at least two distinct nontrivial solutions for $\lambda \in \Lambda_-$,
 99 provided δ is sufficiently small. Furthermore, there is always a one-sided neighborhood Λ_1
 100 of μ_k such that the problem has at least three distinct nontrivial stationary solutions for
 101 $\lambda \in \Lambda_1 \setminus \{\mu_k\}$.

102 It is worth mentioning that “dual” versions of all our results on (3) and (5) mentioned
 103 above hold true if, instead of (4), we assume

$$104 \quad (6) \quad \limsup_{s \rightarrow +\infty} f(x, s) \leq -\bar{f} < 0, \quad \liminf_{s \rightarrow -\infty} f(x, s) \geq \underline{f} > 0$$

105 uniformly for $x \in \bar{\Omega}$.

106 This work is organized as follows. In section 2 we make some preliminaries. In section 3 we
 107 first prove some new global continuation results by utilizing the theory of Conley index. Then
 108 we apply these results to establish a general dynamical bifurcation theorem from infinity for
 109 infinite dynamical systems. In section 4 we use the abstract results to prove some bifurcation
 110 theorems from infinity for nonlinear evolution equations. Finally in section 5, we discuss the
 111 dynamic bifurcation from infinity and multiplicity of stationary solutions for the parabolic
 112 equation mentioned above.

113 **2. Preliminaries.** In this section we make some preliminaries.

114 **2.1. Basic topological notions and results.** Let X be a complete metric space with metric
 115 $d(\cdot, \cdot)$.

Let A and B be nonempty subsets of X . The *distance* $d(A, B)$ between A and B is defined
 as

$$d(A, B) = \inf\{d(x, y) : x \in A, y \in B\},$$

and the *Hausdorff semi-distance* and *Hausdorff distance* of A and B are defined, respectively,
 as

$$d_H(A, B) = \sup_{x \in A} d(x, B), \quad \delta_H(A, B) = \max\{d_H(A, B), d_H(B, A)\}.$$

116 The closure, interior and boundary of A in X are denoted, respectively, by \bar{A} , $\text{int } A$ and
 117 ∂A . Sometimes we also write \bar{A} , $\text{int } A$ and ∂A as \bar{A}^X , $\text{int}_X A$ and $\partial_X A$, respectively, to
 118 emphasize in which space these operations are taken.

119 The ε -neighborhood of A , denoted by $B(A, \varepsilon)$ or $B_X(A, \varepsilon)$, is defined to be the set $\{y \in$
 120 $X : d(y, A) < \varepsilon\}$.

121 A subset U of X is called a *neighborhood* of A , if $\bar{A} \subset \text{int } U$.

122 **Lemma 2.1.** [30] *Let X be a compact metric space, and let A and B be two disjoint closed*
 123 *subsets of X . Then either there exists a subcontinuum C of X such that*

$$124 \quad A \cap C \neq \emptyset \neq B \cap C,$$

126 or $X = X_A \cup X_B$, where X_A and X_B are disjoint compact subsets of X containing A and B ,
 127 respectively.

128 **Lemma 2.2.** ([3, pp. 41]), Let X be a compact metric space. Denote $\mathcal{C}(X)$ the family of
 129 compact subsets of X which is equipped with the Hausdorff metric $\delta_H(\cdot, \cdot)$. Then $\mathcal{C}(X)$ is a
 130 compact metric space.

2.2. Wedge/smash product of pointed spaces. Let (X, x_0) and (Y, y_0) be two pointed spaces. The *wedge product* $(X, x_0) \vee (Y, y_0)$ and *smash product* $(X, x_0) \wedge (Y, y_0)$ are defined, respectively, as follows:

$$(X, x_0) \vee (Y, y_0) = (\mathcal{W}, (x_0, y_0)), \quad (X, x_0) \wedge (Y, y_0) = ((X \times Y)/\mathcal{W}, [\mathcal{W}]),$$

131 where $\mathcal{W} = X \times \{y_0\} \cup \{x_0\} \times Y$.

132 Denote $[(X, x_0)]$ the *homotopy type* of a pointed space (X, x_0) . Since the operations “ \vee ”
 133 and “ \wedge ” preserve homotopy equivalence relations, they can be naturally extended to the
 134 homotopy types of pointed spaces.

Let $\bar{0}$ be the homotopy type of the one-point space $(\{p\}, p)$. Denote Σ^m ($m \geq 0$) the homotopy type of a pointed m -dimensional sphere. Then

$$[(X, x_0)] \vee \bar{0} = [(X, x_0)], \quad \text{and} \quad \Sigma^m \wedge \Sigma^n = \Sigma^{m+n} \quad (\forall m, n \geq 0).$$

135 **2.3. Local semiflows on metric spaces.** For completeness and the reader’s convenience,
 136 let us first collect some fundamental notions and facts on local semiflows.

137 **2.3.1. Local semiflows.** Let X be a complete metric space.

138 A *local semiflow* Φ on X is a continuous mapping from an open set $\mathcal{D}(\Phi) \subset \mathbb{R}_+ \times X$ to X
 139 that enjoys the following properties:

(1) for each $x \in X$, there exists $0 < T_x \leq \infty$, called the *escape time* of $\Phi(t, x)$, such that

$$(t, x) \in \mathcal{D}(\Phi) \iff t \in [0, T_x);$$

(2) $\Phi(0, \cdot) = id_X$, and

$$\Phi(t + s, x) = \Phi(t, \Phi(s, x))$$

140 for all $x \in X$ and $t, s \in \mathbb{R}_+$ with $t + s \leq T_x$.

141 Let Φ be a given local semiflow on X . For simplicity, we usually rewrite $\Phi(t, x)$ as $\Phi(t)x$.

Let $I \subset \mathbb{R}$ be an interval. A *trajectory* (or *solution*) of Φ on I is a continuous mapping $\gamma : I \rightarrow X$ such that

$$\gamma(t) = \Phi(t - s)\gamma(s), \quad \forall t, s \in I, t \geq s.$$

142 A trajectory γ on \mathbb{R} is called a *full trajectory*.

The *orbit* of a trajectory γ on I is the set

$$\text{orb}(\gamma) = \{\gamma(t) : t \in I\}.$$

143 The orbit of a full trajectory is simply called a *full orbit*.

The ω -*limit set* $\omega(\gamma)$ and ω^* -*limit set* of a full trajectory γ are defined as

$$\omega(\gamma) = \{y \in X : \text{there exists } t_n \rightarrow \infty \text{ such that } \gamma(t_n) \rightarrow y\},$$

$$\omega^*(\gamma) = \{y \in X : \text{there exists } t_n \rightarrow -\infty \text{ such that } \gamma(t_n) \rightarrow y\}.$$

Given $U \subset X$, denote $K_\infty(\Phi, U)$ the union of all bounded full orbits in U . In the case where $U = X$, we will simply write

$$K_\infty(\Phi, X) = K_\infty(\Phi).$$

Let $N \subset X$. We say that Φ does not explode in N , if

$$\Phi([0, T_x])x \subset N \implies T_x = \infty.$$

144 **Definition 2.1.** [32] $N \subset X$ is said to be admissible, if for any sequences $x_n \in N$ and
 145 $t_n \rightarrow \infty$ with $\Phi([0, t_n])x_n \subset N$ for all n , the sequence $\Phi(t_n)x_n$ has a convergent subsequence.
 146 N is said to be strongly admissible, if in addition, Φ does not explode in N .

147 **Definition 2.2.** Φ is said to be asymptotically compact on X , if each bounded set $B \subset X$ is
 148 strongly admissible.

149 Let $S \subset X$. S is said to be positively invariant (resp. invariant), if $\Phi(t)S \subset S$ (resp.
 150 $\Phi(t)S = S$) for all $t \geq 0$.

151 A compact invariant set $\mathcal{A} \subset X$ is called an attractor of Φ , if it attracts a neighborhood
 152 U of itself, namely, $\lim_{t \rightarrow +\infty} d_H(\Phi(t)U, \mathcal{A}) = 0$.

153 Let S be a compact invariant set of Φ . An ordered collection $\mathcal{M} = \{M_1, \dots, M_l\}$ of
 154 disjoint compact invariant subsets of S is called a Morse decomposition of S , if for any full
 155 trajectory γ contained in $S \setminus \left(\bigcup_{1 \leq k \leq l} M_k\right)$, there exist i and j with $i < j$ such that

$$156 \quad (1) \quad \omega^*(\gamma) \subset M_j, \quad \omega(\gamma) \subset M_i.$$

157 **Remark 2.1.** A full trajectory satisfying (1) will be referred to as a connecting trajectory
 158 between M_i and M_j .

159 **Remark 2.2.** One may use equivalent definitions of Morse decompositions; see e.g. [32,
 160 Chap. III].

161 **2.4. Conley index.** In this subsection we briefly recall the definition of Conley index. The
 162 interested reader is referred to [7, 25] and [32], etc. for details.

163 Let Φ be a local semiflow on X . Since X may be an infinite dimensional space, we always
 164 assume Φ is asymptotically compact, hence each bounded subset of X is strongly admissible.

165 A compact invariant set S of Φ is said to be isolated, if there is a neighborhood N of
 166 S such that S is the maximal compact invariant set in \bar{N} . Correspondingly, N is called an
 167 isolating neighborhood of S .

168 **Remark 2.3.** Note that we do not require an isolating neighborhood to be bounded, although
 169 the bounded ones are always of particular interest.

170 An important example for isolating neighborhoods is the so called isolating block, which
 171 plays a crucial role in the computation of Conley index.

172 Let $B \subset X$ be a bounded closed set. $x \in \partial B$ is called a strict egress (resp. strict ingress,
 173 bounce-off) point of B , if for every trajectory $\gamma : [-\tau, s] \rightarrow X$ with $\gamma(0) = x$, where $\tau \geq 0$,
 174 $s > 0$, the following properties hold:

(1) there exists $0 < \varepsilon < s$ such that

$$\gamma(t) \notin B \text{ (resp. } \gamma(t) \in \text{int}B, \text{ resp. } \gamma(t) \notin B), \quad \forall t \in (0, \varepsilon);$$

(2) if $\tau > 0$, then there exists $0 < \delta < \tau$ such that

$$\gamma(t) \in \text{int}B \text{ (resp. } \gamma(t) \notin B, \text{ resp. } \gamma(t) \notin B), \quad \forall t \in (-\delta, 0).$$

175 Denote B^e (resp. B^i, B^b) the set of all strict egress (resp. ingress, bounce-off) points of the
 176 closed set B , and set $B^- = B^e \cup B^b$.

177 B is called an *isolating block* [32], if B^- is closed and $\partial B = B^i \cup B^-$.

Let N, E be two closed subsets of X . E is called an *exit set* of N , if (1) E is *N -positively invariant*, that is, for any $x \in E$ and $t \geq 0$,

$$\Phi([0, t])x \subset N \implies \Phi([0, t])x \subset E;$$

178 and (2) for any $x \in N$, if $\Phi(t_1)x \notin N$ for some $t_1 > 0$, then there exists $t_0 \in [0, t_1]$ such that
 179 $\Phi(t_0) \in E$.

180 Let S be a compact isolated invariant set. A pair of bounded closed subsets (N, E) is
 181 called an *index pair* of S , if (1) $N \setminus E$ is an isolating neighborhood of S ; and (2) E is an exit
 182 set of N . We infer from [32] that if B is a bounded isolating block, then (B, B^-) is an index
 183 pair of the maximal compact invariant set $S = K_\infty(\Phi, B)$ in B .

184 **Definition 2.3.** *The homotopy Conley index of S , denoted by $h(\Phi, S)$, is defined to be the*
 185 *homotopy type $[(N/E, [E])]$ of the pointed space $(N/E, [E])$ for any index pair (N, E) of S .*

Remark 2.4. *For convenience, if U is an isolating neighborhood of a compact invariant set S (U need not to be bounded), we also write*

$$h(\Phi, U) = h(\Phi, S),$$

186 *hoping that this will not cause any confusion.*

187 **Example 2.1.** As an example (and also for later use), let us compute the Conley index of an
 188 asymptotically stable equilibrium e (e is an attractor of Φ).

189 Let $L(x)$ be a Lyapunov function of e defined on an open neighborhood U of e which is
 190 strictly decreasing along each trajectory of Φ in U outside e (see e.g. [15, pp. 226] for the
 191 construction of such a function). We may assume $L(e) = 0$ (hence $L(x) > 0$ for $x \in U \setminus \{e\}$).
 192 Take a $\delta > 0$ sufficiently small so that $B = \{x : L(x) \leq \delta\} \subset U$ and is a closed neighborhood
 193 of e . Then one easily sees that B is an isolating block with $B^- = \emptyset$.

We claim that B is contractible. Indeed, set

$$H(s, x) = \begin{cases} \Phi(s/(1-s))x, & x \in B, s \in [0, 1); \\ x, & x \in B, s = 1. \end{cases}$$

194 Then H is a strong deformation retraction.

Now by the definition of Conley index, we have

$$h(\Phi, \{e\}) = [(B/B^-, [B^-])] = \Sigma^0.$$

195 Let S be a compact isolated invariant set of Φ . Denote H_* and H^* the singular homology
 196 and cohomology theories with coefficient group \mathbb{Z} , respectively. Applying H_* and H^* to
 197 $h(\Phi, S)$ one immediately obtains the *homology* and *cohomology Conley indices* of S .

The *Poincaré polynomial* of S , denoted by $p(t, S)$, is the *formal polynomial*

$$p(t, S) = \sum_{q=0}^{\infty} \beta_q t^q$$

with $\beta_q = \text{rank } H_q(h(\Phi, S))$. If S has a Morse decomposition $\mathcal{M} = \{M_1, \dots, M_l\}$, then the following *Morse equation*

$$p(t, M_1) + \dots + p(t, M_l) = p(t, S) + (1+t)Q(t)$$

198 holds for some formal polynomial $Q(t) = \sum_{q=0}^{\infty} d_q t^q$ with $d_q \in \mathbb{Z}_+$.

199 Let us also recall briefly the basic continuation property of the Conley index.

Let Φ_λ ($\lambda \in \Lambda$) be a family of semiflows on X , where Λ is a metric space. We say that Φ_λ *depends on λ continuously*, if $\Phi_\lambda(t)x$ is defined at the point (t, x, λ) , then for any sequence (t_n, x_n, λ_n) converging to (t, x, λ) , $\Phi_{\lambda_n}(t_n)x_n$ is defined as well for all n sufficiently large, furthermore,

$$\Phi_{\lambda_n}(t_n)x_n \rightarrow \Phi_\lambda(t)x \quad \text{as } n \rightarrow \infty.$$

Suppose Φ_λ depends on λ continuously. Define

$$\Pi(t)(x, \lambda) = (\Phi_\lambda(t)x, \lambda), \quad (x, \lambda) \in \mathcal{X} = X \times \Lambda.$$

200 Then Π is a local semiflow on the product space \mathcal{X} , which will be called the *skew-product*
 201 *flow* of the family Φ_λ ($\lambda \in \Lambda$).

202 We say that Φ_λ ($\lambda \in \Lambda$) is λ -*locally uniformly asymptotically compact* (λ -l.u.a.c. in short),
 203 if its skew-product flow Π is asymptotically compact.

204 **Remark 2.5.** *It is trivial to see that if Φ_λ ($\lambda \in \Lambda$) is λ -l.u.a.c., then Λ is necessarily locally*
 205 *compact.*

For convenience, given $K \subset \mathcal{X}$ and $\lambda \in \Lambda$, we will write

$$K[\lambda] = \{x : (x, \lambda) \in K\}.$$

206 $K[\lambda]$ is called the λ -*section* of K . The following continuation result is actually a particular
 207 case of [32, Chap. I, Theorem 12.2].

208 **Theorem 2.1.** *Let Φ_λ ($\lambda \in \Lambda$) be a family of semiflows on X , where Λ is a connected*
 209 *compact metric space. Suppose Φ_λ depends on λ continuously and is λ -l.u.a.c.*

Let K be a compact isolated invariant set of the skew-product flow Π of Φ_λ ($\lambda \in \Lambda$). Then

$$h(\Phi_\lambda, K[\lambda]) \equiv \text{const.}, \quad \lambda \in \Lambda.$$

210 **Proof.** Take a bounded closed isolating neighborhood \mathcal{U} of K in \mathcal{X} . Then for each $\lambda \in \Lambda$,
 211 the λ -section \mathcal{U}_λ of \mathcal{U} is an isolating neighborhood of $K[\lambda]$. By the compactness of K one
 212 can easily verify that $K[\lambda]$ is upper semicontinuous in λ . Consequently \mathcal{U}_λ is also an isolating
 213 neighborhood of $K[\lambda']$ for λ' near λ . The conclusion then directly follows from [32, Chap. I,
 214 Theorem 12.2.]. ■

215 **Remark 2.6.** We emphasize that in the above theorem, we allow $K[\lambda'] = \emptyset$ for some $\lambda' \in \Lambda$.
 216 Note also that when such a case occurs, one necessarily has $h(\Phi_\lambda, K[\lambda]) = \bar{0}$ for all $\lambda \in \Lambda$.

217 **2.5. Sectorial operators.** For the readers' convenience, we finally recall some basic notions
 218 concerning sectorial operators.

Let X be a Banach space. A closed and densely defined linear operator $A : D(A) \subset X \rightarrow X$ is called a *sectorial operator*, if there exist real numbers $\phi \in (0, \pi/2)$, $a \in \mathbb{R}$ and $M \geq 1$ such that the sector

$$S_{a,\phi} = \{\lambda : \phi \leq |\arg(\lambda - a)| \leq \pi, \quad \lambda \neq a\}$$

is contained in the resolvent set of A , moreover,

$$\|(\lambda I - A)^{-1}\| \leq M/|\lambda - a|$$

219 for all $\lambda \in S_{a,\phi}$, where I denotes the identity on X .

Let A be a sectorial operator in X . Denote $\sigma(A)$ the spectral of A . If $\min_{z \in \sigma(A)} \operatorname{Re} z > 0$, then A generates an analytic semigroup $T(t) = e^{-At}$ with

$$\|T(t)\| \leq C e^{-\beta t}, \quad t \geq 0$$

for some $C, \beta > 0$. This allows us to define the fractional powers of A as follows: for each $\alpha > 0$,

$$A^{-\alpha} = \frac{1}{\Gamma(\alpha)} \int_0^\infty t^{\alpha-1} e^{-At} dt,$$

220 where $\Gamma(s) = \int_0^\infty t^{s-1} e^{-t} dt$ is the Gamma function, and let A^α be the inverse of $A^{-\alpha}$ with
 221 $D(A^\alpha) = R(A^{-\alpha})$; see Henry [14, Chap. I] for details. We also assign $A^0 = I$.

Note that in general we may not have $\min_{z \in \sigma(A)} \operatorname{Re} z > 0$. However, one can always find a real number a such that $\min_{z \in \sigma(A_1)} \operatorname{Re} z > 0$, where $A_1 = A + aI$. Hence we can define the fractional powers of A_1 as above. For each $\alpha \geq 0$, denote $X^\alpha = D(A_1^\alpha)$. We equip X^α with the norm $\|\cdot\|_\alpha$ defined as

$$\|u\|_\alpha = \|A_1^\alpha u\|, \quad u \in X^\alpha.$$

222 Then X^α is a Banach space, which is called the *fractional power* of X . It is well known that
 223 the definition of X^α is independent of the choice of the number a , and different choices of a
 224 give equivalent norms on X^α [14, Chap. I].

225 **3. Continuation Theorems and Bifurcation from Infinity of Local Semiflows.** In this
 226 section, we establish some abstract continuation theorems on invariant sets and prove a general
 227 result on bifurcation from infinity in the framework of local semiflows by using Conley index.

Let X be a complete metric space with metric $d(\cdot, \cdot)$, and set

$$\mathcal{X} = X \times \mathbb{R}, \quad \mathcal{X}_\pm = X \times \mathbb{R}_\pm.$$

\mathcal{X} is equipped with the metric defined by

$$\varrho((u_1, \lambda_1), (u_2, \lambda_2)) = d(u_1, u_2) + |\lambda_1 - \lambda_2|, \quad (u_1, \lambda_1), (u_2, \lambda_2) \in \mathcal{X}.$$

228 **3.1. Global continuation theorem.** Let Φ_λ ($\lambda \in \mathbb{R}$) be a family of local semiflows on X .
 229 Henceforth we always assume that Φ_λ depends on λ continuously and is λ -l.u.a.c.
 230 Given $\Lambda \subset \mathbb{R}$ and $U \subset X$, denote

$$231 \quad (1) \quad \mathcal{K}(\Lambda, U) = \overline{\bigcup_{\lambda \in \Lambda} (K_\infty(\Phi_\lambda, U) \times \{\lambda\})}.$$

For simplicity, we will write

$$\mathcal{K}(\Lambda, X) = \mathcal{K}(\Lambda).$$

232
 233 **Remark 3.1.** By the λ -l.u.a.c. property of Φ_λ ($\lambda \in \mathbb{R}$) and the invariance property of
 234 $\mathcal{K}(\Lambda, U)$, one can easily verify that if Λ and U are bounded then $\mathcal{K}(\Lambda, U)$ is compact.

235 **Theorem 3.1.** Let S be a compact isolated invariant set of Φ_0 , and U an isolating neigh-
 236 borhood of S . Denote \mathcal{F}_\pm the family of components of $\mathcal{K}(\mathbb{R}_\pm)$ meeting $S \times \{0\}$.

237 Suppose $h(\Phi_0, S) \neq \bar{0}$. Then there is a $\Gamma \in \mathcal{F}_\pm$ such that either $\Gamma[0] \setminus U \neq \emptyset$, or Γ is
 238 unbounded in the space \mathcal{X}_\pm .

239 *Proof.* We only consider the case of \mathcal{F}_+ . The argument for that of \mathcal{F}_- is parallel.

240 We argue by contradiction and suppose the assertion in the theorem was false. Then each
 241 $\Gamma \in \mathcal{F}_+$ would be bounded in \mathcal{X}_+ . Furthermore, $\Gamma[0] \subset U$ (hence $\Gamma[0] \subset S$).

242 Denote $\mathcal{C}(S)$ the family of all components of S . For each $Z \in \mathcal{C}(S)$, there is a (unique)
 243 $\Gamma_Z \in \mathcal{F}_+$ such that $Z \subset \Gamma_Z[0]$ (note that $\Gamma_Z[0]$ may not be connected). It can be easily seen
 244 that for any $Z_1, Z_2 \in \mathcal{C}(S)$, one has

$$245 \quad (2) \quad \text{either } \Gamma_{Z_1} = \Gamma_{Z_2}, \text{ or } \Gamma_{Z_1} \cap \Gamma_{Z_2} = \emptyset.$$

Let $Z \in \mathcal{C}(S)$. Pick a number δ with $0 < \delta < d(Z, \partial U)$, and let $\mathcal{V}_\delta = B_{\mathcal{X}_+}(\Gamma_Z, \delta)$ be the
 δ -neighborhood of Γ_Z in \mathcal{X}_+ . Set

$$\mathcal{K} = \bar{\mathcal{V}}_\delta \cap \mathcal{K}(\mathbb{R}_+), \quad \mathcal{K}_\delta = \partial_+ \mathcal{V}_\delta \cap \mathcal{K}(\mathbb{R}_+),$$

where $\partial_+ \mathcal{V} = \partial_{\mathcal{X}_+} \mathcal{V}$ denotes the boundary of \mathcal{V} in \mathcal{X}_+ for any $\mathcal{V} \subset \mathcal{X}_+$. Then by the
 boundedness of \mathcal{V}_δ and **Remark 3.1** we easily deduce that both \mathcal{K} and \mathcal{K}_δ are compact. Because
 Γ is a component of $\mathcal{K}(\mathbb{R}_+)$ and $\Gamma_Z \cap \mathcal{K}_\delta = \emptyset$, by virtue of **Lemma 2.1** there exist two disjoint
 closed subsets $\mathcal{K}_1, \mathcal{K}_2$ of \mathcal{K} with $\mathcal{K} = \mathcal{K}_1 \cup \mathcal{K}_2$ such that

$$\Gamma_Z \subset \mathcal{K}_1, \quad \mathcal{K}_\delta \subset \mathcal{K}_2.$$

246 Note that \mathcal{K}_1 is contained in the interior of \mathcal{V}_δ in \mathcal{X}_+ .

Take a number $\delta_Z > 0$ with

$$\delta_Z < \frac{1}{4} \min\{\varrho(\mathcal{K}_1, \mathcal{K}_2), \varrho(\mathcal{K}_1, \partial_+ \mathcal{V}_\delta)\}.$$

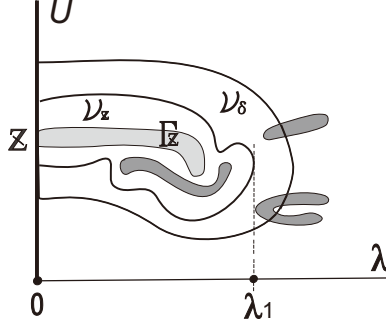
247 Let $\mathcal{V}_Z = B_{\mathcal{X}_+}(\mathcal{K}_1, 2\delta_Z)$. Then by the choice of δ_Z we have

$$248 \quad (3) \quad B_{\mathcal{X}_+}(\partial_+ \mathcal{V}_Z, \delta_Z) \cap \mathcal{K}(\mathbb{R}_+) = \emptyset.$$

249 Let $\lambda_1 = \max\{\lambda : \mathcal{V}_Z[\lambda] \neq \emptyset\}$. Thanks to [Theorem 2.1](#), one deduces that

$$250 \quad (4) \quad h(\Phi_\lambda, \mathcal{V}_Z[\lambda]) \equiv \text{const.}, \quad \lambda \in [0, \lambda_1).$$

251 But $K_\infty(\Phi_\lambda, \bar{\mathcal{V}}_Z[\lambda]) = \emptyset$ if λ is close to λ_1 ; see Fig. 3.1.



252

253

Fig. 3.1: \mathcal{V}_Z is an isolating neighborhood.

By (4) it follows that

$$h(\Phi_\lambda, \mathcal{V}_Z[\lambda]) = \bar{0}, \quad \lambda \in [0, \lambda_1).$$

254 In particular, we have

$$255 \quad (5) \quad h(\Phi_0, \Omega_Z) = \bar{0}, \quad \text{where } \Omega_Z = \mathcal{V}_Z[0].$$

256 Note that $\Omega_Z \subset U$. We also infer from (3) that

$$257 \quad (6) \quad B_X(\partial\Omega_Z, \delta_Z) \cap S = \emptyset.$$

258 (Here $\partial\Omega_Z$ is the boundary of Ω_Z in X .) As S is the maximal compact invariant set of Φ_0 in
259 U , (6) implies that Ω_Z is an isolating neighborhood of Φ_0 .

Since S is compact, there exist a finite number of components Z_1, \dots, Z_l of S such that $S \subset \bigcup_{i=1}^l \Omega_{Z_i}$. Let $W_1 = \Omega_{Z_1}$, and

$$W_k = \Omega_{Z_k} \setminus (\bar{\Omega}_{Z_1} \cup \dots \cup \bar{\Omega}_{Z_{k-1}}), \quad k = 2, \dots, l.$$

260 Then W_k 's are disjoint open sets in X , and

$$261 \quad (7) \quad \partial W_k \subset \bigcup_{i=1}^k \partial\Omega_{Z_i}.$$

As $S \cap \left(\bigcup_{i=1}^l \partial\Omega_{Z_i}\right) = \emptyset$ (see (6)), one finds that

$$S \subset \left(\bigcup_{i=1}^l \Omega_{Z_i}\right) \setminus \left(\bigcup_{i=1}^l \partial\Omega_{Z_i}\right) = \bigcup_{i=1}^l W_i.$$

Set $S_k = S \cap W_k$. We observe that if $w \in S_k$, then by (6),

$$d(w, \partial\Omega_{Z_i}) \geq \delta_{Z_i} \geq \min_{1 \leq i \leq l} \delta_{Z_i} > 0, \quad 1 \leq i \leq l.$$

262 Thus by (7) it holds that

$$263 \quad (8) \quad d(S_k, \partial W_k) > 0,$$

264 which implies that S_k is compact. We also infer from (8) that W_k is an isolating neighborhood
265 of S_k (with respect to Φ_0). We claim that

$$266 \quad (9) \quad h(\Phi_0, S_k) = \bar{0}.$$

Indeed, let $M_k = K_\infty(\Phi_0, \Omega_{Z_k}) \setminus S_k$. Then $M_k \subset \Omega_{Z_k} \setminus W_k$. Therefore by (8) we deduce that

$$d(S_k, M_k) > 0,$$

from which one can easily see that M_k is compact. (5) then asserts that

$$\begin{aligned} \bar{0} = h(\Phi_0, \Omega_{Z_k}) &= h(\Phi_0, K_\infty(\Phi_0, \Omega_{Z_k})) \\ &= h(\Phi_0, S_k \cup M_k) = h(\Phi_0, S_k) \vee h(\Phi_0, M_k). \end{aligned}$$

267 By the basic knowledge in the theory of Conley index (see e.g. [32, pp. 52]) one immediately
268 concludes the validity of (9).

Now since S_k are disjoint isolated invariant sets of Φ_0 and $S = \bigcup_{1 \leq k \leq l} S_k$, we have

$$h(\Phi_0, S) = h(\Phi_0, S_1) \vee \cdots \vee h(\Phi_0, S_l) = \bar{0},$$

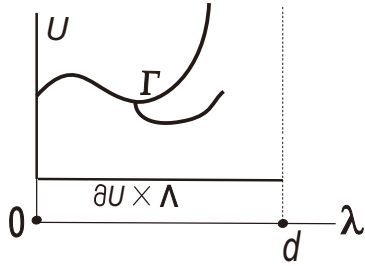
269 which leads to a contradiction. ■

270 **Remark 3.2.** In [40], Ward gave a continuation theorem asserting that $\mathcal{S}_\pm = \bigcup_{\Gamma \in \mathcal{F}_\pm} \Gamma$
271 either meets $(X \setminus U) \times \{0\}$, or is unbounded. **Theorem 3.1** significantly improves this result.

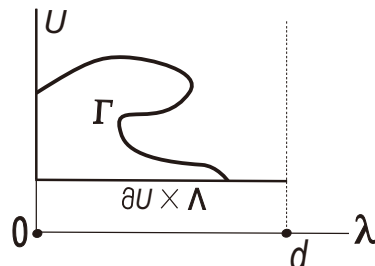
272 **Theorem 3.2.** Let S be an isolated invariant set of Φ_0 with $h(\Phi_0, S) \neq \bar{0}$, and U an isolating
273 neighborhood of S . Let $0 < d \leq \infty$, and denote Λ either the interval $[0, d)$ or the one $(-d, 0]$.
274 Denote \mathcal{F} the family of components of $\mathcal{K}(\Lambda, U)$ meeting $S \times \{0\}$.

275 Then there exists $\Gamma \in \mathcal{F}$ such that one of the alternatives below holds:

- 276 (1) Γ is unbounded; see Fig. 3.2.
277 (2) Γ meeting $\partial U \times \Lambda$; see Fig. 3.3.
278 (3) $\Gamma[\lambda] \neq \emptyset$ for all $\lambda \in \Lambda$; see Fig. 3.4.



279 Fig. 3.2



280 Fig. 3.3

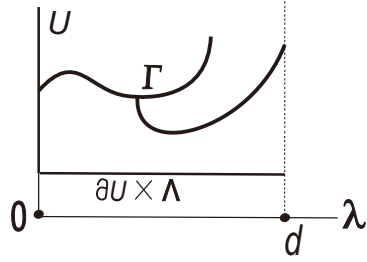
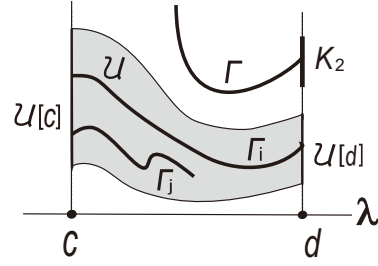


Fig. 3.4


 Fig. 3.5: $K_\infty(\Phi_c) \subset \mathcal{U}[c]$

281

282

283 *Proof.* The proof can be easily obtained by slightly modifying the one of [Theorem 3.1](#).
 284 We omit the details. ■

285 **3.2. An abstract theorem on bifurcation from infinity.** We now establish a new abstract
 286 theorem on dynamic bifurcation from infinity.

287 Let Φ_λ ($\lambda \in \mathbb{R}$) be as in subsection 3.1.

288 **Theorem 3.3.** Let $\Lambda = [c, d]$ be a compact interval. Suppose both $K_\infty(\Phi_c)$ and $K_\infty(\Phi_d)$ are
 289 compact, furthermore,

$$290 \quad (10) \quad h(\Phi_c, K_\infty(\Phi_c)) \neq h(\Phi_d, K_\infty(\Phi_d)).$$

291 Then the set $\mathcal{K}(\Lambda, X)$ has an unbounded component Γ meeting $X \times \{c, d\}$.

Proof. Denote \mathcal{T} the family of connected components of $\mathcal{K}(\Lambda, X)$, and let

$$\mathcal{T}_c = \{\Gamma \in \mathcal{T} : \Gamma[c] \neq \emptyset\}, \quad \mathcal{T}_d = \{\Gamma \in \mathcal{T} : \Gamma[d] \neq \emptyset\}.$$

292 In the following we prove that if every $\Gamma \in \mathcal{T}_c$ is bounded, then there is a $\Gamma \in \mathcal{T}_d$ such that Γ
 293 is unbounded.

294 Let $\mathcal{H} = X \times [c, d]$. Denote $\partial_{\mathcal{H}}\mathcal{V}$ the boundary of \mathcal{V} in \mathcal{H} for any $\mathcal{V} \subset \mathcal{H}$.

Let $\Gamma \in \mathcal{T}_c$. Since Γ is bounded, as in [Remark 3.1](#) one easily deduces by the λ -l.u.a.c. property of Φ_λ that Γ is compact. Take a number $\varepsilon > 0$, and let

$$\mathcal{V}_\varepsilon = \mathcal{B}_{\mathcal{H}}(\Gamma, \varepsilon) := \{(x, \lambda) \in \mathcal{H} : \varrho((x, \lambda), \Gamma) < \varepsilon\}$$

be the ε -neighborhood of Γ in \mathcal{H} . Set

$$\mathcal{C} = \overline{\mathcal{V}_\varepsilon} \cap \mathcal{K}(\Lambda, X), \quad \mathcal{C}_\varepsilon = \partial_{\mathcal{H}}\mathcal{V}_\varepsilon \cap \mathcal{C}.$$

By [Remark 3.1](#) we see that both \mathcal{C} and \mathcal{C}_ε are compact. Since Γ does not intersect any other component of \mathcal{C} , by [Lemma 2.1](#) there exist two disjoint closed subsets \mathcal{C}_1 and \mathcal{C}_2 of \mathcal{C} with $\mathcal{C} = \mathcal{C}_1 \cup \mathcal{C}_2$ such that

$$\Gamma \subset \mathcal{C}_1, \quad \mathcal{C}_\varepsilon \subset \mathcal{C}_2.$$

295 Clearly \mathcal{C}_1 is contained in the interior of \mathcal{V}_ε in \mathcal{H} .

Pick a number $\varepsilon_\Gamma > 0$ with

$$\varepsilon_\Gamma < \frac{1}{4} \min\{\varrho(\mathcal{C}_1, \mathcal{C}_2), \varrho(\mathcal{C}_1, \partial_{\mathcal{H}}\mathcal{V}_{\varepsilon_\Gamma})\}.$$

296 Let $\mathcal{U}_\Gamma = B_{\mathcal{H}}(\mathcal{C}_1, 2\varepsilon_\Gamma)$ be the $2\varepsilon_\Gamma$ -neighborhood of \mathcal{C}_1 in \mathcal{H} . Then by the choice of ε_Γ we see
 297 that $\mathcal{U}_\Gamma \subset \mathcal{V}_\varepsilon$, and moreover,

$$298 \quad (11) \quad B_{\mathcal{H}}(\partial_{\mathcal{H}}\mathcal{U}_\Gamma, \varepsilon_\Gamma) \cap \mathcal{K}(\Lambda, X) = \emptyset.$$

Now we observe that $\mathcal{U} = \{\mathcal{U}_\Gamma[c]\}_{\Gamma \in \mathcal{T}_c}$ forms an open covering of $K_\infty(\Phi_c)$ in X . Thus there exist $\Gamma_1, \dots, \Gamma_n \in \mathcal{T}_c$ such that

$$K_\infty(\Phi_c) \subset \bigcup_{1 \leq i \leq n} \mathcal{U}_{\Gamma_i}[c].$$

299 Let $\mathcal{U} = \bigcup_{1 \leq i \leq n} \mathcal{U}_{\Gamma_i}$. We infer from (11) that \mathcal{U} is an isolating neighborhood of the skew-
 300 product flow Π of $\{\Phi_\lambda\}_{\lambda \in \Lambda}$ in \mathcal{H} with $K_\infty(\Phi_c) \subset \mathcal{U}[c]$; see Fig. 3.5. Therefore by [Theorem 2.1](#)
 301 one concludes that

$$302 \quad (12) \quad h(\Phi_c, K_\infty(\Phi_c)) = h(\Phi_c, \mathcal{U}[c]) = h(\Phi_d, \mathcal{U}[d]) = h(\Phi_d, K_1),$$

303 where $K_1 = K_\infty(\Phi_d, \mathcal{U}[d])$.

For any component Γ of $\mathcal{K}(\Lambda, X)$, by (11) we have $\Gamma \cap \partial_{\mathcal{H}}\mathcal{U}_{\Gamma_i} = \emptyset$ for all $1 \leq i \leq n$. Hence one finds that

$$\text{either } \Gamma \subset \mathcal{U}, \text{ or } \Gamma \cap \overline{\mathcal{U}} = \emptyset.$$

Consequently, for any component C of $K_\infty(\Phi_d)$, we have

$$\text{either } C \subset \mathcal{U}[d], \text{ or } C \cap \overline{\mathcal{U}}[d] = \emptyset.$$

Thus we deduce that $K_\infty(\Phi_d) = K_1 \cup K_2$, where

$$K_2 = \bigcup \{C : C \text{ is a component of } K_\infty(\Phi_d) \text{ with } C \cap \overline{\mathcal{U}}[d] = \emptyset\}.$$

304 As K_1 is isolated with $\mathcal{U}[d]$ being an isolating neighborhood, it is trivial to check that K_2 is
 305 isolated as well. Thereby

$$306 \quad (13) \quad h(\Phi_d, K_\infty(\Phi_d)) = h(\Phi_d, K_1) \vee h(\Phi_d, K_2).$$

This, along with (10) and (12), yields that

$$h(\Phi_d, K_2) \neq \bar{0}.$$

307 Now by virtue of [Theorem 3.1](#), one immediately concludes that there is a $\Gamma \in \mathcal{T}_d$ with
 308 $\Gamma[d] \subset K_2$ such that Γ is unbounded; see Fig. 3.5. ■

309 **3.3. Two examples.** In this subsection we give two simple illustrating examples by con-
 310 sidering ODE systems, which may help the reader have a better understanding to the abstract
 311 results given above.

312 *Example 3.1.* Consider the planar system

$$313 \quad (14) \quad \begin{cases} \dot{x} = x - \lambda x(x^2 + y^2), & x = x(t) \in \mathbb{R}, \\ \dot{y} = y - \lambda y(x^2 + y^2), & y = y(t) \in \mathbb{R}, \end{cases}$$

314 where λ is the bifurcation parameter.

315 Denote Φ_λ the semiflow on $X = \mathbb{R}^2$ generated by the system. Multiplying the first equation
316 in (14) by x and the second one by y , summing the results we obtain that

$$317 \quad (15) \quad \frac{d}{dt}r^2 = 2r^2(1 - \lambda r^2),$$

318 where $r^2 = x^2 + y^2$. Let $\lambda \leq 0$. Then by (15) we have

$$319 \quad (16) \quad \frac{d}{dt}r^2 = 2r^2(1 - \lambda r^2) \geq 2r^2,$$

by which we deduce that $K_\infty(\Phi_\lambda) = \{(0,0)\}$ and is a repeller of the system. Let $B = \overline{B}(2)$, where $B(r)$ denotes the ball in X centered at $(0,0)$ with radius r . By (16) it is clear that B is an isolating block of $K_\infty(\Phi_\lambda)$ with $B^- = \partial B(2)$. Hence

$$h(\Phi_\lambda, K_\infty(\Phi_\lambda)) = [(B/B^-, [B^-])] = \Sigma^2.$$

320 Now assume $\lambda > 0$. By (15) we find that

$$321 \quad (17) \quad \frac{d}{dt}r^2 \leq -r^2$$

as long as $r(t) \geq \sqrt{2/\lambda}$, from which it can be easily seen that the system is dissipative with $K_\infty(\Phi_\lambda)$ being the global attractor. Let $\lambda = 1$. Then we infer from (17) that $B = \overline{B}(2)$ is an isolating block of $K_\infty(\Phi_1)$ with $B^- = \emptyset$. Since B is contractible, one has

$$h(\Phi_1, K_\infty(\Phi_1)) = [(B/B^-, [B^-])] = \Sigma^0.$$

322 Let $\Lambda = [-1, 1]$. Then $h(\Phi_{-1}, K_\infty(\Phi_{-1})) \neq h(\Phi_1, K_\infty(\Phi_1))$. By [Theorem 3.3](#) one immedi-
323 ately concludes that the set $\mathcal{K}(\Lambda, X)$ has an unbounded component Γ meeting $X \times \{\pm 1\}$.

324 One can also discuss the bifurcation phenomena of the system by choosing appropriate
325 isolating neighborhoods of the system and applying [Theorem 3.2](#). For instance, take $U =$
326 $X \setminus B(\frac{1}{2})$. Then for $\lambda \in [0, 1]$, we have by (15) that

$$327 \quad (18) \quad \frac{d}{dt}r^2 = 2r^2(1 - \lambda r^2) > 0, \quad \text{if } (x(t), y(t)) \in \partial U,$$

328 from which one easily deduces that

$$329 \quad (19) \quad K_\infty(\Phi_\lambda, U) \cap \partial U = \emptyset, \quad \forall \lambda \in [0, 1].$$

330 Since $K_\infty(\Phi_\lambda, U) \subset K_\infty(\Phi_\lambda)$ and hence is compact for all λ , by (19) we find that U is an
331 isolating neighborhood of Φ_λ for each $\lambda \in [0, 1]$.

Set $S = K_\infty(\Phi_1, U)$. We infer from the above argument that $S \subset B := \overline{B}(2) \setminus B(\frac{1}{2})$; furthermore, B is an isolating block of S with $B^- = \emptyset$. We have

$$h(\Phi_1, S) = [(B/\emptyset, [\emptyset])] = [(B \cup \{q\}, q)] \neq \overline{0},$$

332 where q is an element with $q \notin B$. By virtue of [Theorem 3.2](#) one concludes that $\mathcal{H}((0, 1], U)$
 333 has a component Γ_U meeting $S \times \{1\}$ such that one of the alternatives (1)-(3) in the theorem
 334 holds true. We claim that Γ_U is unbounded. To see this, we first observe that $K_\infty(\Phi_0, U) = \emptyset$.
 335 Now we argue by contradiction and suppose the contrary. Then one can easily verify that
 336 $\Gamma_U[\lambda] \subset K_\infty(\Phi_\lambda, U)$ for all $\lambda \in (0, 1]$. It follows by [\(19\)](#) that the second alternative (2) in
 337 [Theorem 3.2](#) does not occur. Thus we necessarily have $\Gamma_U[\lambda] \neq \emptyset$ for all $\lambda \in (0, 1]$. But this
 338 and the boundedness of Γ_U then imply that $\Gamma_U[0] \neq \emptyset$, which leads to a contradiction and
 339 proves our claim.

Now let us give a simple observation that justifies our theoretical results obtained above.
 By [\(15\)](#) we see that the circle

$$C_\lambda : r = r_\lambda := 1/\sqrt{\lambda}$$

340 is a closed orbit of the system for each $\lambda > 0$, which depends on λ continuously. Clearly
 341 $r_\lambda \rightarrow \infty$ as $\lambda \rightarrow 0$.

342 It is also worth mentioning that the bifurcating branches Γ and Γ_U given above may be
 343 different. In fact, it is easy to check that for $\lambda \in (0, 1]$, $\Gamma_U[\lambda]$ consists of exactly the closed
 344 orbit C_λ , whereas $\Gamma[\lambda]$ may contain C_λ and the equilibrium $(0, 0)$ and also the connecting
 345 orbits between them.

346 *Example 3.2.* Consider the following non-autonomous scalar equation

$$347 \quad (20) \quad \dot{x} = -(\lambda + h(t))x + e^{-x^2}$$

348 on \mathbb{R} , where $h \in C(\mathbb{R})$ is a T -periodic function ($T > 0$). To have a better understanding of
 349 the dynamics of the equation, as usual we embed the equation into a cocycle system below:

$$350 \quad (21) \quad \dot{x} = -(\lambda + p(t))x + e^{-x^2}, \quad p \in \mathcal{H},$$

351 where $\mathcal{H} = \{h(\tau + \cdot) : \tau \in \mathbb{R}\}$, which is equipped with the topology of uniform convergence
 352 on $[0, T]$ (and hence on \mathbb{R}). It is a basic knowledge that due to the periodicity of h , \mathcal{H} is
 353 homeomorphic to the unit circle (or, one-dimensional sphere) S^1 .

Let $X = \mathbb{R} \times \mathcal{H}$, and denote $\phi_\lambda(t, p)x_0$ the unique solution of [\(21\)](#) with $x(0) = x_0$. Set

$$\Phi_\lambda(t)(x, p) = (\phi_\lambda(t, p)x, \theta_t p), \quad (x, p) \in X,$$

where θ_t is the translation group on \mathcal{H} ,

$$(\theta_t p)(\cdot) = p(t + \cdot), \quad \forall p \in \mathcal{H}, t \in \mathbb{R}.$$

354 Then Φ_λ is a flow on X , called the *skew-product flow* of [\(21\)](#).

355 For the sake of simplicity, we may assume $\max_{\mathbb{R}} |h(t)| \leq 1$. Let $\Lambda = [-2, 2]$. For $\lambda = -2$,
 356 multiplying the equation [\(21\)](#) by x we find that

$$357 \quad (22) \quad \begin{aligned} \frac{1}{2} \frac{d}{dt} x^2 &= (2 - p(t))x^2 + x e^{-x^2} \\ &\geq x^2 - |x| = |x|(|x| - 1), \end{aligned}$$

from which it is clear that for any solution $x(t)$ of (21), if $|x(t_0)| > 1$ for some $t_0 \in \mathbb{R}$ then $|x(t)| > 1$ for all $t \geq t_0$; moreover, $|x(t)| \rightarrow \infty$ as $t \rightarrow +\infty$. It follows that

$$K_\infty(\Phi_{-2}) \subset [-1, 1] \times \mathcal{H}.$$

Let $B_1 = [-2, 2] \times \mathcal{H}$. Making use of (22) it is trivial to check that B_1 is an isolating block of $K_\infty(\Phi_{-2})$ with $B_1^- = \{\pm 2\} \times \mathcal{H}$. Thus

$$h(\Phi_{-2}, K_\infty(\Phi_{-2})) = [(B_1/B_1^-, [B_1^-])].$$

358 Since B_1 is pass-connected, one can easily verify that the quotient space B_1/B_1^- is pass-
 359 connected as well. Hence

360 (23)
$$H_0(h(\Phi_{-2}, K_\infty(\Phi_{-2}))) = H_0((B_1/B_1^-, [B_1^-])) = 0.$$

Now we consider the case where $\lambda = 2$. A fully analogous argument as above applies to show that $K_\infty(\Phi_2) \subset [-1, 1] \times \mathcal{H}$ with $B_2 = [-2, 2] \times \mathcal{H}$ being an isolating block with $B_2^- = \emptyset$. ($K_\infty(\Phi_2)$ is actually the global attractor of Φ_2 .) Thus we have

$$\begin{aligned} h(\Phi_2, K_\infty(\Phi_2)) &= [(B_2/\emptyset, [\emptyset])] = [([[-2, 2] \times \mathcal{H}]/\emptyset, [\emptyset])] \\ &= [(S^1/\emptyset, [\emptyset])] = [(S^1 \cup \{q\}, q)], \end{aligned}$$

361 where q is an element with $q \notin S^1$. Therefore

362 (24)
$$H_0(h(\Phi_2, K_\infty(\Phi_2))) = H_0((S^1 \cup \{q\}, q)) = \mathbb{Z}.$$

363 (23) and (24) indicate that $h(\Phi_{-2}, K_\infty(\Phi_{-2})) \neq h(\Phi_2, K_\infty(\Phi_2))$. Applying **Theorem 3.3**
 364 one immediately concludes that the system Φ_λ undergoes a dynamic bifurcation from infinity
 365 as λ varies in the interval Λ , although we know little about where and how this bifurcation
 366 occurs.

367 **4. Bifurcation from Infinity of Nonlinear Evolution Equations.** In this section we use our
 368 general results in **section 3** to discuss the bifurcation phenomena from infinity of the nonlinear
 369 evolution equation

370 (1)
$$\frac{du}{dt} + Au - \lambda u - f(u, \lambda) = 0$$

371 on a Banach space X , where A is a sectorial operator on X with compact resolvent, $\lambda \in \mathbb{R}$,
 372 and $f(u, \lambda)$ is a locally Lipschitz continuous mapping from $X^\alpha \times \mathbb{R}$ to X for some $0 \leq \alpha < 1$.
 373 Our main goal is to present some more precise descriptions on the dynamic bifurcation from
 374 infinity.

375 Denote $\|\cdot\|$ and $\|\cdot\|_\alpha$ the norms of X and X^α , respectively.

376 **4.1. Existence of unbounded bifurcating branch.** It is well known (see e.g. [14, Theorem
377 3.3.3]) that the Cauchy problem of (1) is well-posed in X^α , that is, for any $u_0 \in X^\alpha$, there
378 exist $T > 0$ and a (unique) continuous function $u : [0, T) \rightarrow X^\alpha$ with $u(0) = u_0$, called
379 the *strong solution* of the problem, such that $u(t) \in D(A)$ and $\frac{d}{dt}u(t)$ exists for $t \in (0, T)$,
380 moreover, the differential equation (1) is satisfied on $(0, T)$.

381 Denote Φ_λ the local semiflow generated by the equation. By the continuity property of
382 f in λ , one can easily verify that Φ_λ depends on λ continuously. Also, by very standard
383 argument (see e.g. [32, Chap. I, Theorem 4.4]), it can be shown that the family Φ_λ ($\lambda \in \mathbb{R}$) is
384 λ -l.u.a.c.

385 We always assume f satisfies the following *sublinear condition*:

386 **(A)** $\lim_{\|u\|_\alpha \rightarrow \infty} \|f(u, \lambda)\| / \|u\|_\alpha = 0$ uniformly on compact λ -intervals.

387 Hence Φ_λ is actually a global semiflow on X^α for each λ .

Definition 4.1. We say that (1) bifurcates from infinity at $\lambda = \mu$ (or, (∞, μ) is a bifurcation
point), if for any $\varepsilon > 0$, there exist $\lambda \in \mathbb{R}$ with $|\lambda - \mu| < \varepsilon$ and a bounded full solution $u_\lambda = u_\lambda(t)$
of (1) such that

$$\|u_\lambda\|_\infty > 1/\varepsilon,$$

388 where $\|u_\lambda\|_\infty = \sup_{t \in \mathbb{R}} \|u_\lambda(t)\|_\alpha$.

Denote $\sigma(A)$ the spectral of A , and write

$$\operatorname{Re} \sigma(A) = \{\operatorname{Re} z : z \in \sigma(A)\}.$$

389

390 **Theorem 4.1.** Let $\mu \in \operatorname{Re} \sigma(A)$. Then (∞, μ) is a bifurcation point of (1). Specifically, for
391 any $c, d \in \mathbb{R}$ with $c < \mu < d$ and $\operatorname{Re} \sigma(A) \cap [c, d] = \{\mu\}$, the set $\mathcal{X}([c, d])$ (see (1) in section 3
392 for the definition) has a component Γ meeting $X^\alpha \times \{c, d\}$ such that for some sequence $\lambda_n \rightarrow \mu$,
393

394 (2) $\sup\{\|x\|_\alpha : x \in \Gamma[\lambda_n]\} \rightarrow \infty$ as $n \rightarrow \infty$.

395 *Proof.* Let us begin with the following linear equation

396 (3)
$$\frac{du}{dt} + Au - \lambda u = 0.$$

397 Let c, d be the numbers given in the theorem. Then if $\lambda = c, d$, the set $\{0\}$ is an isolated
398 invariant set for the semiflow ϕ_λ in X^α generated by (3). By [32] (see Chap. I, Corollary 11.2)
399 there exist two nonnegative integers p and q with $q - p > 0$ such that

400 (4)
$$h(\phi_c, \{0\}) = \Sigma^p, \quad h(\phi_d, \{0\}) = \Sigma^q.$$

401 ($q - p$ is actually the total algebraic multiplicity of all the eigenvalues z of the operator A with
402 $\operatorname{Re} z = \mu$.)

403 Now consider the nonlinear equation

404 (5)
$$\frac{du}{dt} + Au - \lambda u - \nu f(u, \lambda) = 0,$$

where $\nu \in [0, 1]$ is the homotopy parameter. By appropriately modifying the argument in the proof of [32, Chap. II, Theorem 5.1] (see also the proof of [40, Theorem 3.2]), it can be shown that for any $\varepsilon > 0$ with

$$c < \mu - \varepsilon < \mu + \varepsilon < d,$$

405 there exists $R_\varepsilon > 0$ such that for any bounded full solution $u = u(t)$ of (5) with $\lambda \in [c, \mu -$
 406 $\varepsilon] \cup [\mu + \varepsilon, d]$ and $\nu \in [0, 1]$, we have

$$407 \quad (6) \quad \|u\|_\infty < R_\varepsilon.$$

408 Denote ϕ_λ^ν the semiflow generated by (5). By virtue of the continuation property of Conley
 409 index, we conclude that

$$410 \quad (7) \quad \begin{aligned} h(\Phi_\lambda, K_\infty(\Phi_\lambda)) &= h(\phi_\lambda^1, K_\infty(\phi_\lambda^1)) \\ &= h(\phi_\lambda^0, K_\infty(\phi_\lambda^0)) = h(\phi_\lambda, \{0\}) = \Sigma^p \end{aligned}$$

411 for $\lambda \in [c, \mu - \varepsilon]$, and

$$412 \quad (8) \quad \begin{aligned} h(\Phi_\lambda, K_\infty(\Phi_\lambda)) &= h(\phi_\lambda^1, K_\infty(\phi_\lambda^1)) \\ &= h(\phi_\lambda^0, K_\infty(\phi_\lambda^0)) = h(\phi_\lambda, \{0\}) = \Sigma^q \end{aligned}$$

for $\lambda \in [\mu + \varepsilon, d]$. Thanks to Theorem 3.3, one immediately concludes that $\mathcal{K}([c, d])$ has an unbounded connected component Γ meeting $X^\alpha \times \{c, d\}$. On the other hand, (6) implies that for any $\varepsilon > 0$,

$$\Gamma[\lambda] \subset B_{X^\alpha}(R_\varepsilon), \quad \forall \lambda \in [c, \mu - \varepsilon] \cup [\mu + \varepsilon, d],$$

413 where $B_{X^\alpha}(R_\varepsilon)$ denotes the ball in X^α centered at 0 with radius R_ε . Thus there exists a
 414 sequence $\lambda_n \rightarrow \mu$ such that (2) holds true. ■

415 **Remark 4.1.** In Theorem 4.1 one should distinguish two cases of the bifurcation. One is
 416 that $K_\infty(\Phi_\mu)$ is unbounded. When this occurs we say that (1) undergoes a vertical bifurcation
 417 from infinity at $\lambda = \mu$. The other is that $K_\infty(\Phi_\mu)$ is bounded, in which case we deduce that
 418 there is a sequence $\lambda_n \rightarrow \mu$ ($\lambda_n \neq \mu$ for all n) such that $\Gamma[\lambda_n]$ is unbounded, where Γ is the
 419 connected bifurcating branch given in the theorem. Note that both cases may occur. This can
 420 be seen from the following two simple examples.

421 *Example 4.1.* Consider the linear equation

$$422 \quad (9) \quad \dot{u} + u = \lambda u, \quad u = u(t) \in \mathbb{R},$$

423 where $\lambda \in \mathbb{R}$ is the bifurcation parameter. Then we can see that $\mu = 1$ is a bifurcation value,
 424 at which each constant function $u(t) = c$ ($c \in \mathbb{R}$) is a bounded full solution of the equation.
 425 Hence the equation undergoes a vertical bifurcation from infinity at $\lambda = 1$.

426 It is also interesting to note that for each $\lambda \neq 1$, the equation has no bounded full solutions
 427 other than the trivial one.

428 *Example 4.2.* Consider the non-homogenous equation

$$429 \quad (10) \quad \dot{u} + u = \lambda u + 1, \quad u \in \mathbb{R},$$

430 where $\lambda \in \mathbb{R}$ is the bifurcation parameter. Again $\mu = 1$ is a bifurcation value, at which each
 431 solution of (10) is given by $u = t + c$ ($c \in \mathbb{R}$). Clearly $K_\infty(\Phi_\mu) = \emptyset$.

On the other hand, if we let $[c, d] = [0, 2]$, then by [Theorem 4.1](#) we see that $\mathcal{K}([0, 2])$
 has an unbounded connected component Γ in the space $\mathbb{R} \times [0, 2]$ with $\Gamma \cap (\mathbb{R} \times \{0, 2\}) \neq \emptyset$.
 Actually, for $\lambda \neq 1$, the unique bounded full solution of the equation is the stationary one
 $u_\lambda(t) = (1 - \lambda)^{-1}$. Hence

$$\Gamma = \{(u_\lambda, \lambda) : 0 \leq \lambda < 1\}$$

432 is a component of $\mathcal{K}([0, 2])$ fulfilling all the requirements in the theorem.

433 **4.2. Further results on dynamic bifurcation from infinity.** We infer from [Theorem 4.1](#)
 434 that there is a sequence $\lambda_n \rightarrow \mu$ such that for each $\lambda = \lambda_n$, (1) has a bounded full solution
 435 $u_n = u_n(t)$ with $\|u_n\|_\infty \rightarrow \infty$. In what follows we give another result on the bifurcation of
 436 the equation from infinity, which seems to be more precise in some aspects.

437 **Theorem 4.2.** *Assume f satisfies the sublinear condition (A) in [Theorem 4.1](#). Let $\mu \in$
 438 $\text{Re } \sigma(A)$. Then one of the following alternatives holds.*

439 (1) *There is a sequence u_n of bounded full solutions of (1) at $\lambda = \mu$ such that $\lim_{n \rightarrow \infty} \|u_n\|_\infty =$
 440 ∞ .*

441 (2) *There is a one-sided neighborhood Λ_1 of μ such that for each $\lambda \in \Lambda_1 \setminus \{\mu\}$, (1) has two
 442 distinct bounded full solutions u_λ and v_λ such that*

$$443 \quad (11) \quad \lim_{\lambda \rightarrow \mu} \|u_\lambda\|_\infty = \infty,$$

444 *whereas $\|v_\lambda\|_\infty$ remains bounded on the λ -interval Λ_1 .*

445 (3) *There is a two-sided neighborhood Λ of μ such that for each $\lambda \in \Lambda \setminus \{\mu\}$, the equation (1)
 446 has a bounded full solution u_λ satisfying (11).*

447 *Proof.* If (1) holds true then we are done. Thus we assume the contrary, and hence S_μ is
 448 a bounded set, where (and below) $S_\lambda = K_\infty(\Phi_\lambda)$.

449 Take two numbers $c, d \in \mathbb{R}$ as in [Theorem 4.1](#). Since the number ε in (7) and (8) is
 450 arbitrary, we infer from (7) and (8) that

$$451 \quad (12) \quad h(\Phi_\lambda, S_\lambda) = \Sigma^p \ (\lambda \in [c, \mu]), \quad h(\Phi_\lambda, S_\lambda) = \Sigma^q \ (\lambda \in (\mu, d])$$

452 for some nonnegative integers p and q with $p < q$.

453 Pick a bounded closed isolating neighborhood U of S_μ . Choose a $\delta > 0$ sufficiently small
 454 so that U is an isolating neighborhood of Φ_λ for all $\lambda \in \Lambda = [\mu - \delta, \mu + \delta]$. Then

$$455 \quad (13) \quad h(\Phi_\lambda, U) \equiv \text{const.}$$

456 Two possibilities may occur.

457 Case 1) $h(\Phi_\mu, S_\mu) \neq \bar{0}$. In such a case we show that the second assertion (2) holds true.
 It is obvious that

$$\text{either } h(\Phi_\mu, S_\mu) \neq \Sigma^p, \quad \text{or } h(\Phi_\mu, S_\mu) \neq \Sigma^q.$$

458 Let us first consider the case where $h(\Phi_\mu, S_\mu) \neq \Sigma^p$. By (12) we have

459 (14)
$$h(\Phi_\lambda, S_\lambda) \neq h(\Phi_\mu, S_\mu), \quad \lambda \in [c, \mu).$$

460 We claim that

461 (15)
$$S_\lambda \setminus U \neq \emptyset, \quad \forall \lambda \in \Lambda_- := [\mu - \delta, \mu).$$

Indeed, if $S_\lambda \subset U$ for some $\lambda \in \Lambda_-$, then by (12) and (13) one finds that

$$h(\Phi_\mu, S_\mu) = h(\Phi_\mu, U) = h(\Phi_\lambda, U) = h(\Phi_\lambda, S_\lambda) = \Sigma^p,$$

462 which leads to a contradiction.

463 For each $\lambda \in \Lambda_-$, pick an $x_\lambda \in S_\lambda \setminus U$. Let u_λ be a full trajectory of Φ_λ contained in S_λ
 464 with $u_\lambda(0) = x_\lambda$. We show that u_λ fulfills (11).

465 Suppose the contrary. Then there would exist a sequence $\lambda_n \rightarrow \mu$ ($\lambda_n \neq \mu$) such that
 466 the sequence $u_n = u_{\lambda_n}$ is uniformly bounded on \mathbb{R} . By very standard argument it can be
 467 shown that u_n has a subsequence converging to a bounded full trajectory u_0 of Φ_μ uniformly
 468 on any compact interval of \mathbb{R} . u_0 is necessarily contained in S_μ . On the other hand, since
 469 $u_n(0) = x_{\lambda_n} \notin U$, we deduce that $u_0(0) \notin \text{int } U$. This leads to a contradiction.

470 Now assume that $h(\Phi_\mu, S_\mu) \neq \Sigma^q$. Then by a fully analogous argument as above, one
 471 concludes that for each $\lambda \in \Lambda_+ = (\mu, \mu + \delta]$, the equation has a bounded full solution u_λ
 472 satisfying (11).

Since $h(\Phi_\mu, S_\mu) \neq \bar{0}$, by (13) we have

$$h(\Phi_\lambda, U) = h(\Phi_\mu, U) = h(\Phi_\mu, S_\mu) \neq \bar{0}, \quad \lambda \in \Lambda.$$

473 It follows that $K_\infty(\Phi_\lambda, U) \neq \emptyset$. For each $\lambda \in \Lambda$, pick a full solution v_λ in $K_\infty(\Phi_\lambda, U)$. Then
 474 $\|v_\lambda\|_\infty$ remains bounded on Λ .

Case 2) $h(\Phi_\mu, S_\mu) = \bar{0}$. In this case, we have

$$\Sigma^p \neq h(\Phi_\mu, S_\mu) \neq \Sigma^q.$$

475 The same argument as in Case 1) applies to show that for each $\lambda \in \Lambda_- \cup \Lambda_+$, the equation
 476 has a bounded full solution u_λ satisfying (11). Hence the assertion (3) holds. ■

477 **5. Dynamic Bifurcation and Multiplicity for Parabolic Equations.** In this section we
 478 consider the following boundary value problem:

479 (1)
$$\begin{cases} u_t - \Delta u = \lambda u + f(x, u), & x \in \Omega; \\ u(x, t) = 0, & x \in \partial\Omega, \end{cases}$$

480 where Ω is a bounded domain in \mathbb{R}^n , $\lambda \in \mathbb{R}$, and $f \in C^1(\bar{\Omega} \times \mathbb{R})$.

Let $H = L^2(\Omega)$ and $V = H_0^1(\Omega)$. By (\cdot, \cdot) and $|\cdot|$ we denote the usual inner product and
 norm on H , respectively. The norm $\|\cdot\|$ on V is defined by

$$\|u\| = \left(\int_\Omega |\nabla u|^2 dx \right)^{1/2}, \quad u \in V.$$

Denote A the operator $-\Delta$ associated with the homogenous Dirichlet boundary condition. A is a sectorial operator and has a compact resolvent. Denote

$$0 < \mu_1 < \mu_2 < \cdots < \mu_k < \cdots$$

481 the eigenvalues of A .

482 We may convert (1) into an abstract equation on V :

$$483 \quad (2) \quad u_t + Au = \lambda u + \tilde{f}(u), \quad u = u(t) \in V,$$

where $\tilde{f}(u)$ is the Nemitski operator from V to H given by

$$\tilde{f}(u)(x) = f(x, u(x)), \quad u \in V.$$

484 If we assume that

$$485 \quad (3) \quad f(x, s) = o(|s|) \quad \text{as } |s| \rightarrow \infty$$

486 uniformly with respect to $x \in \bar{\Omega}$, then one can trivially verify that the Nemitski operator \tilde{f}
487 in (2) satisfies the sublinear condition **(A)** in section 4. Thus applying the abstract results in
488 section 4, one can immediately obtain some interesting information on the bifurcation of the
489 equation. For instance, we have

490 **Theorem 5.1.** *Let μ_k be an eigenvalue of A . Then one of the following alternatives holds.*

(1) *There is a sequence u_n of bounded full solutions of (2) at $\lambda = \mu_k$ such that*

$$\lim_{n \rightarrow \infty} \|u_n\|_{\infty} = \infty.$$

491 (2) *There is a one-sided neighborhood Λ_1 of μ_k such that for $\lambda \in \Lambda_1 \setminus \{\mu_k\}$, the equation (2)
492 has at least two distinct bounded full solutions u_{λ} and v_{λ} such that*

$$493 \quad (4) \quad \lim_{\lambda \rightarrow \mu_k} \|u_{\lambda}\|_{\infty} = \infty,$$

494 whereas $\|v_{\lambda}\|$ remains bounded on Λ_1 .

495 (3) *There is a two-sided neighborhood Λ of μ_k such that for each $\lambda \in \Lambda \setminus \{\mu_k\}$, (2) has at least
496 one bounded full solution u_{λ} satisfying (4).*

497 In this present work, we are basically interested in a particular but very important case,
498 namely, the case where f satisfies the Landesman-Laser type condition (4) in section 1. We will
499 give some precise descriptions on the bifurcation of the equation and discuss the multiplicity
500 of stationary solutions of the equation.

501 Henceforth we always assume

502 **(H)** f satisfies the Landesman-Laser type condition (4) in section 1.

Denote Φ_{λ} the semiflow associated with (2), namely, for each $u_0 \in X^{\alpha}$,

$$u(t) = \Phi_{\lambda}(t)u_0$$

503 is the solution of the equation on \mathbb{R}_+ with initial value $u(0) = u_0$.

504 **5.1. Preliminaries.** Let us begin with a fundamental result on f .

Given a function w on Ω , we use w_{\pm} to denote the positive and negative parts of w , respectively,

$$w_{\pm} = \max\{\pm w(x), 0\}, \quad x \in \Omega.$$

505 Then $w = w_+ - w_-$. We have

Lemma 5.1. *For any $R, \varepsilon > 0$, there exists $s_0 > 0$ such that*

$$\int_{\Omega} f(x, v + sw) w dx \geq \int_{\Omega} (\bar{f} w_+ + \underline{f} w_-) dx - \varepsilon$$

506 for all $s \geq s_0$, $v \in \bar{B}_H(R)$ and $w \in \bar{B}_H(1)$, where $B_H(r)$ denotes the ball in H centered at 0
507 with radius r .

508 *Proof.* This is a slightly modified version of [17, Lemma 6.7]. Here we give the details of
509 the proof for completeness and the reader's convenience.

Let

$$I = \int_{\Omega} f(x, v + sw) w dx - \int_{\Omega} (\bar{f} w_+ + \underline{f} w_-) dx.$$

Since $w = w_+ - w_-$, we can rewrite I as $I_+ - I_-$, where

$$I_+ = \int_{\Omega} (f(x, v + sw) - \bar{f}) w_+ dx, \quad I_- = \int_{\Omega} (f(x, v + sw) + \underline{f}) w_- dx.$$

510 In what follows, let us estimate I_+ for $v \in \bar{B}_H(R)$ and $w \in \bar{B}_H(1)$.

We observe that

$$R^2 \geq \int_{\Omega} |v|^2 dx \geq \int_{\{|v| \geq \sigma\}} |v|^2 dx \geq \sigma^2 |\{|v| \geq \sigma\}|,$$

511 from which it can be easily seen that $|\{|v| \geq \sigma\}| \rightarrow 0$ as $\sigma \rightarrow +\infty$ uniformly with respect to
512 $v \in \bar{B}_H(R)$. (Here and below $|E|$ denotes the Lebesgue measure for any measurable subset E
513 of \mathbb{R}^n .) Therefore there exists $\sigma > 0$ such that

$$514 \quad (5) \quad |\{|v| \geq \sigma\}|^{1/2} < \delta := \varepsilon/8 \|f\| (|\Omega| + 1), \quad v \in \bar{B}_H(R),$$

515 where $\|f\| = \sup_{x \in \bar{\Omega}, s \in \mathbb{R}} |f(x, s)|$.

For each $v \in \bar{B}_H(R)$ and $w \in \bar{B}_H(1)$, let

$$D = D_{v,w} := \{|v| < \sigma\} \cap \{w_+ > \delta\}.$$

Then $\Omega = D \cup \{|v| \geq \sigma\} \cup \{w_+ \leq \delta\}$. Hence

$$\begin{aligned} I_+ &\geq \int_D (f(x, v + sw) - \bar{f}) w_+ dx - \int_{\{|v| \geq \sigma\}} |f(x, v + sw) - \bar{f}| w_+ dx \\ &\quad - \int_{\{w_+ \leq \delta\}} |f(x, v + sw) - \bar{f}| w_+ dx \\ &\geq \int_D (f(x, v + sw) - \bar{f}) w_+ dx - 2\|f\| \left(\int_{\{|v| \geq \sigma\}} w_+ dx + \int_{\{w_+ \leq \delta\}} w_+ dx \right). \end{aligned}$$

Note that

$$\begin{aligned} \int_{\{|v| \geq \sigma\}} w_+ dx &\leq \left(\int_{\{|v| \geq \sigma\}} w_+^2 dx \right)^{1/2} |\{|v| \geq \sigma\}|^{1/2} \\ &\leq (\text{by (5)}) \leq |w| \delta \leq \delta. \end{aligned}$$

It is obvious that

$$\int_{\{w_+ \leq \delta\}} w_+ dx \leq |\Omega| \delta.$$

516 Thereby

$$\begin{aligned} 517 \quad (6) \quad I_+ &\geq \int_D (f(x, v + sw) - \bar{f}) w_+ dx - 2 \|f\| (|\Omega| + 1) \delta \\ &= \int_D (f(x, v + sw) - \bar{f}) w_+ dx - \frac{\varepsilon}{4}. \end{aligned}$$

Since $z + s\eta \rightarrow +\infty$ (as $s \rightarrow +\infty$) uniformly for $z \in [-\sigma, \sigma]$ and $\eta \geq \delta$, there exists $s_1 > 0$ (independent of v and w) such that if $s > s_1$ then

$$f(x, z + s\eta) - \bar{f} \geq -\frac{\varepsilon}{4|\Omega|^{1/2}}, \quad \forall z \in [-\sigma, \sigma], \eta \geq \delta.$$

Now suppose that $s > s_1$. Then by the definition of D , we have (note that $w = w_+$ on D)

$$\begin{aligned} \int_D (f(x, v + sw) - \bar{f}) w_+ dx &\geq -\frac{\varepsilon}{4|\Omega|^{1/2}} \int_D w_+ dx \\ &\geq -\frac{\varepsilon}{4|\Omega|^{1/2}} |D|^{1/2} (\int_D |w|^2 dx)^{1/2} \geq -\frac{\varepsilon}{4}. \end{aligned}$$

Thus by (6) we see that

$$I_+ \geq \int_D (f(x, v + sw) - \bar{f}) w_+ dx - \frac{\varepsilon}{4} > -\frac{\varepsilon}{2}.$$

518 Similarly it can be shown that there exists $s_2 > 0$ (independent of v and w) such that
519 $I_- < \frac{\varepsilon}{2}$, provided $s > s_2$.

Set $s_0 = \max\{s_1, s_2\}$. Then if $s > s_0$, we have

$$I \geq I_+ - I_- > -\frac{\varepsilon}{2} - \frac{\varepsilon}{2} = -\varepsilon$$

520 for all $v \in \bar{B}_H(R)$ and $w \in \bar{B}_H(1)$. This completes the proof of the lemma. ■

521 Now we prove some basic facts concerning the dynamical behavior of the equation (2).

Let $L = A - \mu_k$, where μ_k is an eigenvalue of A . The space H can be decomposed into the orthogonal direct sum of its subspaces H^- , H^0 and H^+ corresponding to the negative, zero and positive eigenvalues of L , respectively. Note that both H^- and H^0 are of finite-dimensional. Denote P^σ ($\sigma \in \{0, \pm\}$) the projection from H to H^σ . Set

$$V^\sigma = V \cap H^\sigma, \quad \sigma \in \{0, \pm\}.$$

By the finite dimensionality of H^- and H^0 , one finds that V^- and V^0 coincide with H^- and H^0 , respectively. We also have

$$V = V^- \oplus V^0 \oplus V^+.$$

Lemma 5.2. *Assume $\lambda \leq \mu_k + \eta$, where $\eta = (\mu_{k+1} - \mu_k)/2$. Then there exists $\rho_0 > 0$ (independent of λ) such that for any solution $u = u(t)$ of (2) on \mathbb{R}_+ ,*

$$\|u^+(t)\|^2 \leq \|u_0^+\|^2 e^{-\eta t} + \rho_0^2(1 - e^{-\eta t}), \quad \forall t \geq 0.$$

522 Here $u^+ = P^+u$.

Proof. Taking the inner product of the equation with Au^+ in H , it yields

$$\begin{aligned} \frac{1}{2} \frac{d}{dt} \|u^+\|^2 + |Au^+|^2 &= \lambda \|u^+\|^2 + (Au^+, \tilde{f}(u)) \\ &\leq \lambda \|u^+\|^2 + \varepsilon |Au^+|^2 + C_\varepsilon \end{aligned}$$

523 for any $\varepsilon > 0$, where C_ε is a positive constant depending only upon ε and the upper bound of
524 $|\tilde{f}(s)|$. Hence

$$525 \quad (7) \quad \frac{1}{2} \frac{d}{dt} \|u^+\|^2 + (1 - \varepsilon) |Au^+|^2 \leq \lambda \|u^+\|^2 + C_\varepsilon.$$

526 Note that $|Au^+|^2 \geq \mu_{k+1} \|u^+\|^2$. Therefore by (7) we have

$$527 \quad (8) \quad \frac{1}{2} \frac{d}{dt} \|u^+\|^2 \leq -((1 - \varepsilon)\mu_{k+1} - \lambda) \|u^+\|^2 + C_\varepsilon.$$

Fix an $\varepsilon > 0$ sufficiently small so that $(1 - \varepsilon)\mu_{k+1} > \mu_k + \frac{3}{2}\eta$. Then for $\lambda \leq \mu_k + \eta$, one has

$$(1 - \varepsilon)\mu_{k+1} - \lambda > \left(\mu_k + \frac{3}{2}\eta \right) - (\mu_k + \eta) = \eta/2.$$

528 Now the conclusion follows from (8) and the classical Gronwall Lemma. ■

Denote

$$\Xi_\rho = \{v \in V : \|P^+v\| \leq \rho\}, \quad \rho > 0.$$

529 As a direct consequence of Lemma 5.2, we have

Corollary 5.1. *Assume $\lambda \leq \mu_k + \eta$. Then*

$$K_\infty(\Phi_\lambda) \subset \Xi_{\rho_0}.$$

530 Furthermore, Ξ_ρ is positively invariant under Φ_λ for any $\rho > \rho_0$.

Set $W = V^- \oplus V^0$, and let

$$P_W = P^- + P^0$$

531 be the projection from V to W . Given $0 \leq a < b \leq \infty$ and $\rho > 0$, denote

$$532 \quad (9) \quad \Xi_\rho[a, b] = \{u \in \Xi_\rho : a \leq |P_W u| \leq b\}.$$

533 **Lemma 5.3.** *Let η and ρ_0 be as in Lemma 5.2, and $\rho > \rho_0$. Then there exist $R_0, c_0 > 0$
534 such that the following assertions hold.*

535 (1) If $\lambda \in [\mu_k, \mu_k + \eta]$, then for any solution $u(t)$ of the equation (2) in $\Xi_\rho[R_0, \infty]$, we have

$$536 \quad (10) \quad \frac{d}{dt}|w(t)|^2 \geq c_0|w(t)|,$$

537 where $w(t) = P_W u(t)$.

538 (2) For any $R > R_0$, there exists $0 < \theta \leq \eta$ such that if $\lambda \in [\mu_k - \theta, \mu_k]$, then (10) holds true
539 for any solution $u(t)$ of (2) in $\Xi_\rho[R_0, R]$.

540 *Proof.* Let $\lambda \in [\mu_k - \eta, \mu_k + \eta]$, and $u = u(t)$ a solution of (2) in Ξ_ρ . Taking the inner
541 product of (2) with $w = P_W u$ in H , it yields

$$542 \quad (11) \quad \frac{1}{2} \frac{d}{dt}|w|^2 + \|w\|^2 = \lambda|w|^2 + (\tilde{f}(u), w).$$

543 Because $\|w\|^2 \leq \mu_k|w|^2$, by (11) we have

$$544 \quad (12) \quad \frac{1}{2} \frac{d}{dt}|w|^2 \geq (\lambda - \mu_k)|w|^2 + (\tilde{f}(u), w).$$

545 Let us first estimate the last term in (12).

546 As the norm $\|\cdot\|_{L^1(\Omega)}$ of $L^1(\Omega)$ and that of $H = L^2(\Omega)$ are equivalent on W , one easily
547 sees that

$$548 \quad (13) \quad \min\{\|v\|_{L^1(\Omega)} : v \in W, |v| = 1\} := m > 0.$$

549 Pick a number $\delta > 0$ with $\delta \leq \min\{\bar{f}, \underline{f}\}$. By virtue of Lemma 5.1 there exists $s_0 > 0$
550 (depending only upon ρ) such that if $s \geq s_0$, then

$$551 \quad (14) \quad (\tilde{f}(h + sv), v) = \int_{\Omega} f(x, h + sv)v \, dx \geq \int_{\Omega} (\bar{f}v_+ + \underline{f}v_-) \, dx - \frac{1}{2}m\delta$$

552 for all $h \in \bar{B}_H(\rho)$ and $v \in \bar{B}_H(1)$.

Now we rewrite

$$w = sv, \quad \text{where } s = |w|.$$

Then $|v| = 1$. Suppose $s \geq s_0$. Noticing that $\|u^+\| \leq \rho$, by (14) one finds that

$$(\tilde{f}(u), w) = s(f(x, u^+ + sv), v) \geq s \left(\int_{\Omega} (\bar{f}v_+ + \underline{f}v_-) \, dx - \frac{1}{2}m\delta \right).$$

Observing that

$$\begin{aligned} & \int_{\Omega} (\bar{f}v_+ + \underline{f}v_-) \, dx - \frac{1}{2}m\delta \\ & \geq \delta \int_{\Omega} |v| \, dx - \frac{1}{2}m\delta \geq (\text{by (13)}) \geq \frac{1}{2}m\delta, \end{aligned}$$

553 we conclude that

$$554 \quad (15) \quad (\tilde{f}(u), w) \geq \frac{1}{2}m\delta s = \frac{1}{2}m\delta|w|.$$

555 Now we combine (15) and (12) together to obtain that

$$556 \quad (16) \quad \frac{d}{dt}|w(t)|^2 \geq 2(\lambda - \mu_k)|w|^2 + m\delta|w(t)|$$

557 as long as $|w(t)| \geq s_0$.

Set $R_0 = s_0$, $c_0 = m\delta/2$. Assume $\lambda \in [\mu_k, \mu_k + \eta]$. Then $\lambda - \mu_k \geq 0$, and we infer from (16) that

$$\frac{d}{dt}|w(t)|^2 \geq m\delta|w(t)| > c_0|w(t)|$$

558 at any point t where $|w(t)| \geq R_0$. Hence the assertion (1) holds.

Now assume $\lambda < \mu_k$. Let $R > R_0$. Choose a $\theta > 0$ with $\theta R^2 < m\delta s_0/4$. Then if $\lambda \in [\mu_k - \theta, \mu_k)$, for any solution $u(t)$ of (2) in $\Xi_\rho[R_0, R]$, by (16) we conclude that

$$\begin{aligned} \frac{d}{dt}|w(t)|^2 &\geq -2|\lambda - \mu_k|R^2 + m\delta|w(t)| \\ &\geq c_0|w(t)| + (c_0|w(t)| - 2\theta R^2) \\ &\geq c_0|w(t)| + (c_0 s_0 - 2\theta R^2) \geq c_0|w(t)|, \end{aligned}$$

559 which justifies the second assertion (2). ■

560 **5.2. Dynamic bifurcation from infinity.** We are now ready to discuss the bifurcation of
561 the equation (2) near $\lambda = \mu_k$.

562 Let Φ_λ be the semiflow generated by (2). First, as a consequence of Lemma 5.3 we have
563 the following basic fact.

564 **Proposition 5.1.** *Assume the hypothesis (H). Then $K_\infty(\Phi_\lambda)$ is uniformly bounded in V for*
565 $\lambda \in [\mu_k, \mu_k + \eta]$, and

$$566 \quad (17) \quad h(\Phi_{\mu_k}, K_\infty(\Phi_{\mu_k})) = h(\Phi_\lambda, K_\infty(\Phi_\lambda)) = \Sigma^{p+r},$$

567 where p is the sum of the multiplicities of the eigenvalues μ_i ($0 \leq i \leq k-1$) of A , and r the
568 multiplicity of μ_k .

569 *Proof.* Let η and ρ_0 be the numbers given in Lemma 5.2. Fix a number $\rho > \rho_0$. Then
570 there exist $R_0, c_0 > 0$ such that the first assertion (1) in Lemma 5.3 holds true, by which one
571 easily deduces that

$$572 \quad (18) \quad K_\infty(\Phi_\lambda) \subset \Xi_\rho[0, R_0], \quad \forall \lambda \in [\mu_k, \mu_k + \eta].$$

573 On the other hand, as in (12) in section 4 it can be shown that

$$574 \quad (19) \quad h(\Phi_\lambda, K_\infty(\Phi_\lambda)) = \begin{cases} \Sigma^{p+r}, & \lambda \in (\mu_k, \mu_{k+1}); \\ \Sigma^p, & \lambda \in (\mu_{k-1}, \mu_k). \end{cases}$$

575 By (18) and the continuation property of Conley index we immediately conclude that

$$576 \quad h(\Phi_{\mu_k}, K_\infty(\Phi_{\mu_k})) = \Sigma^{p+r}. \quad \blacksquare$$

577 Now we state and prove the main result in this subsection on the dynamic bifurcation
578 from infinity of the equation near each eigenvalue μ_k .

579 **Theorem 5.2.** *Assume the hypothesis **(H)**. Then $S_\lambda := K_\infty(\Phi_\lambda)$ is nonvoid for all $\lambda \in \mathbb{R}$,
580 and there exists $\delta > 0$ such that the following assertions hold.*

581 (1) *For each $\lambda \in \Lambda_- := [\mu_k - \delta, \mu_k)$, S_λ has a Morse decomposition $\mathcal{M} = \{M_\lambda^\infty, M_\lambda^1\}$.
582 Furthermore, there is at least one connecting trajectory γ between M_λ^1 and M_λ^∞ .*

583 (2) *M_λ^1 remains uniformly bounded on Λ_- , whereas*

$$584 \quad (20) \quad \lim_{\lambda \rightarrow \mu_k^-} \min_{v \in M_\lambda^\infty} \|v\| = \infty.$$

(3) *Each of the sets \mathcal{K}^1 and \mathcal{K}^∞ has a component Γ with $\Gamma[\lambda] \neq \emptyset$ for all $\lambda \in \Lambda_-$, where*

$$\mathcal{K}^1 = \overline{\bigcup_{\lambda \in \Lambda_-} (M_\lambda^1 \times \{\lambda\})}, \quad \mathcal{K}^\infty = \overline{\bigcup_{\lambda \in \Lambda_-} (M_\lambda^\infty \times \{\lambda\})}.$$

585 *Proof.* (i) We infer from the proof of **Theorem 4.1** (see (6) in section 4) and **Proposition 5.1**
586 that S_λ is a compact subset of V for all $\lambda \in \mathbb{R}$. Since the Conley index of S_λ is nontrivial (see
587 (17) and (19)), one concludes that $S_\lambda \neq \emptyset$.

(ii) **Corollary 5.1** asserts that

$$S_\lambda \subset \Xi_{\rho_0}, \quad \forall \lambda \leq \mu_k + \eta,$$

588 where ρ_0 is the number in **Lemma 5.2**. Fix a $\rho > \rho_0$, and let R_0 and c_0 be the numbers given
589 by **Lemma 5.3**. Pick a bounded isolating neighborhood N_1 of S_{μ_k} with

$$590 \quad (21) \quad \Xi_\rho[0, R_0] \subset N_1.$$

591 Then one can restrict $\delta > 0$ sufficiently small so that N_1 is also an isolating neighborhood of
592 Φ_λ for all $\lambda \in \Lambda := [\mu_k - \delta, \mu_k + \delta]$. Hence

$$593 \quad (22) \quad h(\Phi_\lambda, M_\lambda^1) \equiv \text{const.}, \quad \lambda \in \Lambda,$$

594 where $M_\lambda^1 = K_\infty(\Phi_\lambda, N_1)$. Further by **Proposition 5.1** we deduce that

$$595 \quad (23) \quad h(\Phi_\lambda, M_\lambda^1) = h(\Phi_{\mu_k}, M_{\mu_k}^1) = h(\Phi_{\mu_k}, S_{\mu_k}) = \Sigma^{p+r}, \quad \lambda \in \Lambda.$$

596 It is clear that $M_\lambda^1 \subset S_\lambda \subset \Xi_\rho$. Therefore

$$597 \quad (24) \quad M_\lambda^1 \subset N_1 \cap \Xi_\rho := \tilde{N}_1, \quad \lambda \in \Lambda.$$

598 Because \tilde{N}_1 is bounded, one can find a number $R_1 > 0$ such that

$$599 \quad (25) \quad \tilde{N}_1 \subset \Xi_\rho[0, R_1/2].$$

600 By **Lemma 5.3** (2), there exists $\theta_1 > 0$ such that if $\lambda \in [\mu_k - \theta_1, \mu_k)$, then for any solution
601 $u(t)$ of (2) in $\Xi_\rho[R_0, R_1]$, one has

$$602 \quad (26) \quad \frac{d}{dt}|w(t)|^2 \geq c_0|w(t)| \geq c_0R_0 > 0,$$

603 where $w(t) = P_W u(t)$. We may assume $\delta \leq \theta_1$. Let $\lambda \in \Lambda_- := [\mu_k - \delta, \mu_k)$. Then for any
 604 bounded full solution $u(t)$ of (2) in Ξ_ρ with $u(t_0) \in \Xi_\rho[R_0, R_1]$ for some t_0 , by (26) one easily
 605 deduces that there exists $T > 0$ such that

$$606 \quad (27) \quad u(t) \in \Xi_\rho[0, R_0] \quad (t < -T), \quad \text{and} \quad u(t) \in \Xi_\rho[R_1, \infty] \quad (t > T).$$

607 Combining (24), (25) and (27) together, it yields that

$$608 \quad (28) \quad M_\lambda^1 \subset \Xi_\rho[0, R_0], \quad \forall \lambda \in \Lambda_-.$$

609 As M_λ^1 is the maximal compact invariant set of Φ_λ in N_1 , (21) and (28) imply that M_λ^1 is the
 610 maximal compact invariant set in $\Xi_\rho[0, R_0]$.

Set

$$M_\lambda^\infty = K_\infty(\Phi_\lambda, \Xi_\rho[R_1, \infty]).$$

611 Then $M_\lambda^\infty \subset K_\infty(\Phi_\lambda) = S_\lambda$. We prove that $\mathcal{M} = \{M_\lambda^\infty, M_\lambda^1\}$ forms a Morse decomposition
 612 of S_λ . For this purpose, let us first show that if $u = u(t)$ is a full solution in $S_\lambda \setminus (M_\lambda^1 \cup M_\lambda^\infty)$,
 613 then

$$614 \quad (29) \quad \omega^*(u) \subset M_\lambda^1, \quad \omega(u) \subset M_\lambda^\infty.$$

615 Indeed, let u be such a solution. Then since $S_\lambda \subset \Xi_\rho$ and M_λ^1 and M_λ^∞ are the maximal
 616 compact invariant sets in $\Xi_\rho[0, R_0]$ and $\Xi_\rho[R_1, \infty]$, respectively, there exists $t_0 \in \mathbb{R}$ such that
 617 $u(t_0) \in \Xi_\rho[R_0, R_1]$. Hence (29) directly follows from (27).

Now we check that $M_\lambda^\infty \neq \emptyset$. Thus \mathcal{M} is a Morse decomposition of S_λ . Suppose the
 contrary. Then by (29) we find that $S_\lambda = M_\lambda^1$. Hence

$$h(\Phi_\lambda, M_\lambda^1) = h(\Phi_\lambda, S_\lambda) = (\text{by (19)}) = \Sigma^p,$$

618 which contradicts (23).

619 To complete the proof of (1), there remains to check the existence of a connecting trajectory
 620 between M_λ^1 and M_λ^∞ . To this end, we consider the Morse equation of \mathcal{M} :

$$621 \quad (30) \quad p(t, M_\lambda^1) + p(t, M_\lambda^\infty) = p(t, S_\lambda) + (1+t)Q(t).$$

Recalling that $h(\Phi_\lambda, M_\lambda^1) = \Sigma^{p+r}$ and $h(\Phi_\lambda, S_\lambda) = \Sigma^p$, we have

$$p(t, M_\lambda^1) = t^{p+r}, \quad p(t, S_\lambda) = t^p.$$

622 Thus (30) reads

$$623 \quad (31) \quad t^{p+r} + p(t, M_\lambda^\infty) = t^p + (1+t)Q(t),$$

624 which implies that $Q(t) \neq 0$. By the basic knowledge in the Morse theory of invariant sets (see
 625 [32, Chap. III, Theorem 3.5]), one immediately concludes that there is at least one connecting
 626 trajectory between M_λ^1 and M_λ^∞ .

627 We also infer from (31) that $p(t, M_\lambda^\infty) \neq 0$. Consequently

$$628 \quad (32) \quad h(\Phi_\lambda, M_\lambda^\infty) \neq \bar{0}, \quad \forall \lambda \in \Lambda_-.$$

629 (iii) Clearly M_λ^1 remains uniformly bounded on Λ .

630 For any $R > R_1$, by Lemma 5.3 there exists $0 < \theta < \theta_1$ such that when $\lambda \in [\mu_k - \theta, \mu_k)$,
631 the differential inequality (26) holds true for any solution $u(t)$ of (2) in $\Xi_\rho[R_0, 2R]$. Using this
632 basic fact, it can be easily seen that if $\lambda \in [\mu_k - \theta, \mu_k)$, any bounded full solution in $\Xi_\rho[R_1, \infty]$
633 is necessarily contained in $\Xi_\rho[R, \infty]$. Hence

$$634 \quad (33) \quad M_\lambda^\infty \subset \Xi_\rho[R, \infty],$$

635 which implies what we desired in (20) and completes the proof of (2).

636 (iv) Finally, let us verify the validity of (3).

Let $U = \Xi_\rho[R_1, \infty]$. Then $\partial U = C_1 \cup C_2$, where

$$C_1 = \{v : \|v^+\| = \rho, |w| \geq R_1\}, \quad C_2 = \{v : \|v^+\| \leq \rho, |w| = R_1\}.$$

Here $v^+ = P^+v$, and $w = P_Wv$. Let $\lambda \in \Lambda_-$. By the choice of ρ and Lemma 5.2, we see that
 $M_\lambda^\infty \cap C_1 = \emptyset$. Fix an $R > R_1$. Then we infer from the above argument in (iii) that one can
restrict $\delta > 0$ to be sufficiently small so that (33) holds. Consequently $M_\lambda^\infty \cap C_2 = \emptyset$. Thus

$$M_\lambda^\infty \cap \partial U = \emptyset,$$

637 namely, U is an isolating neighborhood of M_λ^∞ .

638 Because $h(\Phi_{\mu_k - \delta}, U) \neq \bar{0}$ (by (32)), \mathcal{K}^∞ has a connected component Γ with $\Gamma[\mu_k - \delta] \neq \emptyset$
639 such that one of the alternatives (1)-(3) in Theorem 3.2 holds true. As $\Gamma[\lambda] \subset \text{int } U$ for all
640 $\lambda \in \Lambda_-$, we conclude that either Γ is unbounded, or $\Gamma[\mu_k] \neq \emptyset$. Because $\Gamma[\lambda]$ is uniformly
641 bounded on $[\mu_k - \delta, \mu_k - \varepsilon]$ for any $\varepsilon \in (0, \delta)$, in any case we deduce that $\Gamma[\lambda] \neq \emptyset$ for all
642 $\lambda \in \Lambda_-$.

643 The argument for \mathcal{K}^1 is similar. We omit the details. ■

5.3. Bifurcation and multiplicity of stationary solutions. We now turn to the static
bifurcation and multiplicity of stationary solutions of (2). Since the equation has a natural
Lyapunov function $J(u)$ defined by

$$J(u) = \frac{1}{2}(\|u\|^2 - \lambda|u|^2) - \int_\Omega F(x, u)dx, \quad u \in V,$$

644 where $F(x, s) = \int_0^s f(x, t)dt$, this problem can be treated in the framework of dynamical
645 systems.

646 **Theorem 5.3.** *Assume the hypothesis (H). Let $\delta > 0$ be the same as in Theorem 5.2. Then*

647 (1) Φ_λ has at least one equilibrium e_λ for all $\lambda \in \mathbb{R}$;

648 (2) there exists $\delta > 0$ such that Φ_λ has at least two distinct equilibria e_λ^∞ and e_λ^c for each
649 $\lambda \in \Lambda_- = [\mu_k - \delta, \mu_k)$, and

$$650 \quad (34) \quad \lim_{\lambda \rightarrow \mu_k^-} \|e_\lambda^\infty\| = \infty,$$

651 whereas e_λ^c remains bounded on Λ_- ; and

652 (3) there is an open dense subset \mathcal{D} of \mathbb{R} such that for each $\lambda \in \Lambda_- \cap \mathcal{D}$, Φ_λ has at least three
653 distinct equilibria.

654 *Proof.* (1) Since each nonempty compact invariant set contains at least one stationary
655 solution, the conclusion (1) directly follows from [Theorem 5.2](#).

656 (2) Let N_1 be the isolating neighborhood of S_{μ_k} given in the proof of [Theorem 5.2](#), and
657 let $\mathcal{M} = \{M_\lambda^\infty, M_\lambda^1\}$ be the Morse decomposition of S_λ for $\lambda \in \Lambda_-$. Then $M_\lambda^1 \subset N_1$. By (25)
658 we have

$$659 \quad (35) \quad N_1 \cap \Xi_\rho = \tilde{N}_1 \subset \Xi_\rho[0, R_1/2], \quad \lambda \in \Lambda_-.$$

660 As $M_\lambda^\infty \subset \Xi_\rho[R_1, \infty]$, by (35) we find that $M_\lambda^\infty \cap N_1 = \emptyset$. Pick two stationary solutions e_λ
661 and e_λ^∞ from M_λ^1 and M_λ^∞ , respectively. Then e_λ and e_λ^∞ fulfill the requirements in (2).

662 (3) By slightly modifying the proof of [[35](#), Theorem 2.1], it can be shown that there is
663 an open dense subset \mathcal{D} of \mathbb{R} such that all the equilibria of Φ_λ are hyperbolic if $\lambda \in \mathcal{D}$. Now
664 assume $\lambda \in \Lambda_- \cap \mathcal{D}$. We show that there is another equilibrium $z_\lambda^\infty \in M_\lambda^\infty$ with $z_\lambda^\infty \neq e_\lambda^\infty$.
665 Consequently Φ_λ has at least three distinct equilibria.

We argue by contradiction and suppose M_λ^∞ consists of exactly one hyperbolic stationary
solution e_λ^∞ . Then $p(t, M_\lambda^\infty) = t^m$ for some $m \geq 0$. Accordingly the Morse equation ([30](#))
reads

$$t^{p+r} + t^m = t^p + (1+t)Q(t).$$

666 But this is impossible for any formal polynomial $Q(t)$ with coefficients in \mathbb{Z}_+ , as the sum of
667 the coefficients of the left-hand side does not equal that of the right-hand side.

668 The proof of the theorem is finished. ■

669 **Remark 5.1.** We infer from the above argument that for each $\lambda \in \Lambda_- \cap \mathcal{D}$, Φ_λ has at least
670 two distinct equilibria outside the domain N_1 .

671 Finally, we pay some special attention to the particular case where

672 **(F)** $f(x, s) = o(|s|)$ as $|s| \rightarrow 0$ uniformly for $x \in \bar{\Omega}$.

673 We prove some new multiplicity results on stationary solutions for the equation (2) near each
674 eigenvalue μ_k . The main results are summarized in the following theorem.

675 **Theorem 5.4.** Assume f satisfies the hypotheses **(H)** and **(F)**. Denote $W_{loc}^c(0)$ the local
676 center manifold of Φ_{μ_k} at the trivial equilibrium 0, and let ϕ be the restriction of Φ_{μ_k} on
677 $W_{loc}^c(0)$. Suppose 0 is an isolated equilibrium of Φ_{μ_k} . Then there exists $\delta > 0$ such that one
678 of the following assertions holds:

679 (1) 0 is an attractor of ϕ . In this case, the system Φ_λ has at least two distinct nontrivial
680 equilibria e_λ^c and e_λ^∞ for $\lambda \in \Lambda_- = [\mu_k - \delta, \mu_k)$, whereas it has at least three distinct ones e_λ^1 ,
681 e_λ^2 and e_λ^c for $\lambda \in \Lambda_+ = (\mu_k, \mu_k + \delta]$.

682 (2) 0 is a repeller of ϕ (i.e., an attractor of the inverse flow ϕ^{-1}). In this case, Φ_λ has at
683 least three distinct nontrivial equilibria e_λ^1 , e_λ^2 and e_λ^∞ for each $\lambda \in \Lambda_-$.

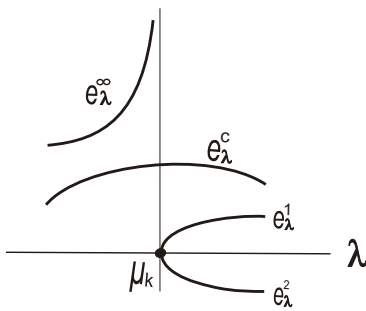
684 (3) 0 is neither an attractor nor a repeller of ϕ . When this occurs, Φ_λ has at least three
 685 nontrivial equilibria e_λ^1 , e_λ^c and e_λ^∞ for $\lambda \in \Lambda_-$, whereas it has at least two distinct ones e_λ^1
 686 and e_λ^c for $\lambda \in \Lambda_+$.

687 Furthermore, we have

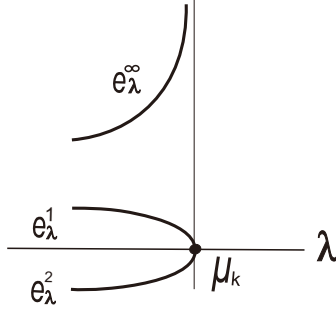
$$688 \quad (36) \quad \lim_{\lambda \rightarrow \mu_k} \|e_\lambda^\infty\| = \infty, \quad \lim_{\lambda \rightarrow \mu_k} \|e_\lambda^i\| = 0 \quad (i = 1, 2),$$

689 and

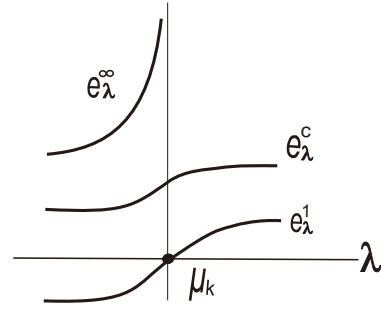
$$690 \quad (37) \quad 0 < \liminf_{\lambda \rightarrow \mu_k} \|e_\lambda^c\| \leq \limsup_{\lambda \rightarrow \mu_k} \|e_\lambda^c\| < \infty.$$



691
692 Fig. 5.1



693 Fig. 5.2



694 Fig. 5.3

693 *Proof.* In the following argument, we always assume that $\delta > 0$ is sufficiently small so that
 694 the conclusions in [Theorem 5.2](#) and [Theorem 5.3](#) are valid.

695 (1) The case where 0 is an attractor of ϕ .

696 Let N_1 be the isolating neighborhood of S_{μ_k} in the proof of [Theorem 5.3](#). Then by
 697 [Theorem 5.3](#), for each $\lambda \in \Lambda_-$ the system Φ_λ always has an equilibrium e_λ^∞ outside N_1
 698 satisfying the first equation in (36).

Pick a number $\beta > 0$ and an isolating neighborhood N_0 of 0 with

$$N_0 \subset B_V(\beta) \subset N_1.$$

We may restrict δ so that both N_0 and N_1 are isolating neighborhoods of Φ_λ for all $\lambda \in \Lambda :=$
 $[\mu_k - \delta, \mu_k + \delta]$. Let

$$K_\lambda^i = K_\infty(\Phi_\lambda, N_i), \quad i = 0, 1.$$

699 Then for each i ,

$$700 \quad (38) \quad h(\Phi_\lambda, K_\lambda^i) \equiv \text{const.}, \quad \lambda \in \Lambda.$$

701 It is trivial to check that

$$702 \quad (39) \quad d_H(K_\lambda^0, \{0\}) \rightarrow 0 \quad \text{as } \lambda \rightarrow \mu_k.$$

It can be assumed that N_0 is sufficiently small so that the product formula of Conley index given in [32, Chap. II, Theorem 3.1] holds true, hence

$$h(\Phi_{\mu_k}, \{0\}) = \Sigma^p \wedge h(\phi, \{0\}).$$

We infer from Example 2.1 that $h(\phi, \{0\}) = \Sigma^0$. Therefore

$$h(\Phi_{\mu_k}, \{0\}) = \Sigma^p \wedge \Sigma^0 = \Sigma^p.$$

703 It then follows by (38) that

$$704 \quad (40) \quad h(\Phi_\lambda, K_\lambda^0) = h(\Phi_{\mu_k}, K_{\mu_k}^0) = h(\Phi_{\mu_k}, \{0\}) = \Sigma^p, \quad \lambda \in \Lambda.$$

705 By (38) and Proposition 5.1 we also deduce that

$$706 \quad (41) \quad h(\Phi_\lambda, K_\lambda^1) = h(\Phi_{\mu_k}, K_{\mu_k}^1) = h(\Phi_{\mu_k}, S_{\mu_k}) = \Sigma^{p+r}, \quad \lambda \in \Lambda.$$

707 Thus we see that $K_\lambda^1 \neq K_\lambda^0$. As K_λ^0 is the maximal invariant set in N_0 , one concludes that

$$708 \quad (42) \quad K_\lambda^1 \setminus N_0 \neq \emptyset, \quad \lambda \in \Lambda.$$

709 For each $\lambda \in \Lambda$, pick a $v_\lambda \in K_\lambda^1 \setminus N_0$. Let $u_\lambda(t)$ be a bounded full trajectory of Φ_λ in K_λ^1
710 with $u_\lambda(0) = v_\lambda$. We claim that if δ is small enough then

$$711 \quad (43) \quad \text{either } \omega(u_\lambda) \setminus N_0 \neq \emptyset, \quad \text{or } \omega^*(u_\lambda) \setminus N_0 \neq \emptyset.$$

712 Indeed, if this was false, there would exist a sequence $\lambda_n \rightarrow \mu_k$ (as $n \rightarrow \infty$) such that both
713 $\omega(u_n)$ and $\omega^*(u_n)$ are contained in N_0 and hence in $K_{\lambda_n}^0$, where $u_n = u_{\lambda_n}$. Thus by (39) we
714 deduce that

$$715 \quad (44) \quad \lim_{n \rightarrow \infty} \max_{v \in \omega(u_n)} |J(v)| = 0 = \lim_{n \rightarrow \infty} \max_{v \in \omega^*(u_n)} |J(v)|.$$

Set

$$\Gamma_n = \overline{\text{orb}(u_n)} = \text{orb}(u_n) \cup \omega(u_n) \cup \omega^*(u_n).$$

Then

$$\min_{v \in \Gamma_n} J(v) = \min_{v \in \omega(u_n)} J(v), \quad \max_{v \in \Gamma_n} J(v) = \max_{v \in \omega^*(u_n)} J(v).$$

716 It follows by (44) that

$$717 \quad (45) \quad \max_{v \in \Gamma_n} |J(v)| \rightarrow 0 \quad \text{as } n \rightarrow \infty.$$

718 On the other hand, since $\Gamma_n \subset K_{\lambda_n}^1 \subset N_1$ and Φ_λ is λ -l.u.a.c., it is easy to verify that
719 $\bigcup_{\lambda_n \in \Lambda} \Gamma_n$ is precompact. Hence by Lemma 2.2 it can be assumed that Γ_n converges to a
720 nonempty compact invariant set K of Φ_{μ_k} (in the sense of Hausdorff distance). Noticing that
721 $\Gamma_n \cap K_{\lambda_n}^0 \neq \emptyset$, by (39) we find that $0 \in K$. Because each Γ_n is connected, K is connected as
722 well. (45) implies that $J(v) \equiv 0$ on K . Thereby each point in K is an equilibrium of Φ_{μ_k} . As

$$723 \quad (46) \quad u_n(0) \in \Gamma_n \setminus N_0$$

724 for all n , we deduce that $K \setminus \text{int } N_0 \neq \emptyset$. Further by the connectedness of K one concludes that
 725 $K \cap \partial V \neq \emptyset$ for any small neighborhood V of 0 , which leads to a contradiction and completes
 726 the proof of our claim.

727 In view of (43), for each $\lambda \in \Lambda$ we can pick an equilibrium e_λ^c of Φ_λ with

$$728 \quad (47) \quad e_\lambda^c \in (\omega(u_\lambda) \cup \omega^*(u_\lambda)) \setminus N_0 \subset N_1 \setminus N_0.$$

729 Hence if $\lambda \in \Lambda_-$, the system Φ_λ has at least two distinct nontrivial equilibria e_λ^c and e_λ^∞ .

730 We infer from the attractor bifurcation theory (see e.g. Ma and Wang [20, Theorem 4.3],
 731 [19, Theorem 6.1] or Li and Wang [18, Theorem 4.2]) that K_λ^0 contains at least two distinct
 732 equilibrium points e_λ^1 and e_λ^2 for $\lambda \in \Lambda_+$, provided δ is sufficiently small. By (47) one concludes
 733 that Φ_λ has at least three distinct nontrivial equilibria for $\lambda \in \Lambda_+$.

734 (2) 0 is a repeller of ϕ . In this case, as in (1), by applying the attractor bifurcation theory
 735 we deduce that K_λ^0 contains at least two distinct equilibria e_λ^1 and e_λ^2 for $\lambda \in \Lambda_-$. Since Φ_λ
 736 has a nontrivial equilibrium e_λ^∞ outside N_1 for each $\lambda \in \Lambda_-$, it has at least three distinct ones
 737 for $\lambda \in \Lambda_-$.

738 (3) Finally, let us consider the case where 0 is neither an attractor nor a repeller of ϕ .
 739 By Li and Wang [18, Theorem 4.4] we deduce that the system bifurcates at each side of μ_k a
 740 nonempty compact invariant set $M_\lambda \subset N_0$ with $0 \notin M_\lambda$ and

$$741 \quad (48) \quad d_H(M_\lambda, \{0\}) \rightarrow 0 \quad \text{as } \lambda \rightarrow \mu_k.$$

742 M_λ contains at least one nontrivial equilibrium e_λ^1 .

743 We show that

$$744 \quad (49) \quad h(\Phi_{\mu_k}, \{0\}) \neq \Sigma^{p+r},$$

745 which fact will yield another equilibrium $e_\lambda^c \in N_1 \setminus N_0$ at both sides of μ_k .

746 Consider the local center-unstable manifold $W_{loc}^{cu}(0)$ of Φ_{μ_k} at 0 . Denote ψ the restriction
 747 of Φ_{μ_k} on $W_{loc}^{cu}(0)$. Then

$$748 \quad (50) \quad h(\Phi_{\mu_k}, \{0\}) = h(\psi, \{0\}).$$

749 Thus to prove (49), it suffices to check that

$$750 \quad (51) \quad H_*(h(\psi, \{0\})) \neq H_*(\Sigma^{p+r}).$$

We argue by contradiction and suppose the contrary. Then

$$H_{p+r}(h(\psi, \{0\})) = H_{p+r}(\Sigma^{p+r}) = \mathbb{Z}.$$

Therefore by the Poincaré-Lefschetz duality theory of the Conley index (see McCord [24,
 Theorem 2.1] and Mrozek and Szrednicki [26, pp. 164]),

$$H^0(h(\psi^{-1}, \{0\})) = H_{p+r}(h(\psi, \{0\})) = \mathbb{Z}.$$

On the other hand, pick a pass-connected isolating block $B \subset W_{loc}^{cu}(0)$ of $S_0 = \{0\}$ with respect to the inverse flow ψ^{-1} . (Such an isolating block is always available due to [8, Theorem 1.5].) Since S_0 is not an attractor of ψ^{-1} (note that S_0 is not an attractor of ϕ^{-1} on $W_{loc}^c(0)$), we necessarily have

$$B^- \neq \emptyset.$$

Thus by the basic knowledge in the theory of algebraic topology, one easily deduces that $H^0(B, B^-) = 0$. Consequently

$$H^0(h(\psi^{-1}, \{0\})) = H^0(B, B^-) = 0,$$

751 which leads to a contradiction and justifies the validity of (49).

Recall that (see (41))

$$h(\Phi_\lambda, K_\lambda^1) = \Sigma^{p+r}, \quad \lambda \in \Lambda = [\mu_k - \delta, \mu_k + \delta].$$

Noticing that

$$h(\Phi_\lambda, K_\lambda^0) = h(\Phi_{\mu_k}, K_{\mu_k}^0) = h(\Phi_{\mu_k}, \{0\}) \neq \Sigma^{p+r}, \quad \forall \lambda \in \Lambda,$$

we conclude that $K_\lambda^1 \neq K_\lambda^0$. As K_λ^0 is the maximal invariant set in N_0 , one finds that

$$K_\lambda^1 \setminus N_0 \neq \emptyset, \quad \lambda \in \Lambda.$$

752 We are now in a quite similar situation as in (42). Repeating the same argument below
753 (42), it can be easily shown that the system has an equilibrium e_λ^c in $N_1 \setminus N_0$.

754 In conclusion, there are at least two distinct nontrivial equilibria in N_1 for $\lambda \in \Lambda \setminus \{\mu_k\}$.
755 Because Φ_λ has an equilibrium e_λ^∞ outside N_1 for $\lambda \in \Lambda_-$, the system has at least three
756 distinct nontrivial equilibria as $\lambda \in \Lambda_-$. This completes the proof of (3).

757 The second equation in (36) follows from (39) and (48). (37) is a direct consequence of
758 the choice that $e_\lambda^c \in N_1 \setminus N_0$. ■

759 **Remark 5.2.** *It is interesting to note that there is always a one-sided neighborhood Λ_1 of*
760 *μ_k such that the equation has at least three distinct nontrivial stationary solutions for $\lambda \in \Lambda_1$.*

761 **Remark 5.3.** *Dual versions of all the results in this section hold true if, instead of (H), we*
762 *assume that (6) in section 1 is fulfilled.*

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